Study of $B$ meson production in $p + Pb$ collisions at $\sqrt{s_{NN}} = 5.02 TeV$ using exclusive hadronic decays

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Abstract: The production cross sections of the $B^+, B^0$, and $B_s^0$ mesons, and of their charge conjugates, are measured via exclusive hadronic decays in $p+Pb$ collisions at the center-of-mass energy $\sqrt{s_{NN}} = 5.02 TeV$ with the CMS detector at the CERN LHC. The data set used for this analysis corresponds to an integrated luminosity of 34.6 nb$^{-1}$. The production cross sections are measured in the transverse momentum range between 10 and 60 GeV/c. No significant modification is observed compared to proton-proton perturbative QCD calculations scaled by the number of incoherent nucleon-nucleon collisions. These results provide a baseline for the study of in-medium $b$ quark energy loss in $Pb+Pb$ collisions.

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Study of $B$ Meson Production in $p + \text{Pb}$ Collisions at $\sqrt{s_{NN}} = 5.02$ TeV Using Exclusive Hadronic Decays

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The production cross sections of the $B^+$, $B^0$, and $B^0_s$ mesons, and of their charge conjugates, are measured via exclusive hadronic decays in $p + \text{Pb}$ collisions at the center-of-mass energy $\sqrt{s_{NN}} = 5.02$ TeV with the CMS detector at the CERN LHC. The data set used for this analysis corresponds to an integrated luminosity of $34.6 \text{ nb}^{-1}$. The production cross sections are measured in the transverse momentum range between 10 and 60 GeV/$c$. No significant modification is observed compared to proton-proton perturbative QCD calculations scaled by the number of incoherent nucleon-nucleon collisions. These results provide a baseline for the study of in-medium $b$ quark energy loss in $\text{Pb} + \text{Pb}$ collisions.

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Relativistic heavy ion collisions allow the study of quantum chromodynamics (QCD) at very high temperature and density. Under such extreme conditions, a strongly interacting state consisting of deconfined quarks and gluons, the quark-gluon plasma (QGP) [1,2], is predicted by lattice QCD calculations [3]. Hard-scattered partons are expected to lose energy as they traverse the QGP via elastic collisions and medium-induced gluon radiation. The resulting reduction of the measured yield of hadrons, compared to expectations based on proton-proton ($pp$) data, is often referred to as “jet quenching” [4,5]. The flavor dependence of jet quenching is one of the most important testing grounds for energy loss models [6–10]. However, other phenomena can affect the yield of heavy-flavor particles, independently of the presence of a deconfined partonic medium. For instance, modifications of the parton distribution functions (PDFs) in the nucleus with respect to nucleon PDFs [11–13] could change the production rate. Therefore, a complete understanding of the interactions of heavy quarks in the deconfined medium formed in heavy ion collisions requires a thorough knowledge of their production in proton- (or deuteron-) nucleus, $p(d) + A$, collisions.

Currently, published data for heavy-flavor production in $p(d) + A$ exist for open charm both at RHIC, in $d + \text{Au}$ collisions at $\sqrt{s_{NN}} = 200$ GeV, and at the LHC in $p + \text{Pb}$ collisions at $\sqrt{s_{NN}} = 5.02$ TeV. At RHIC, the STAR Collaboration measured the charm spectra in the rapidity interval $|y| < 1$ from direct reconstruction of $D^0$ meson and from indirect electron and positron measurements of charm semileptonic decays [14]. The measured yields were found to be consistent within the uncertainties with the hypothesis of binary scaling (no modification with respect to nucleon-nucleon ($NN$) cross section scaled by the number of incoherent $NN$ binary collisions). However, the PHENIX Collaboration measured a significant enhancement of the production of heavy-flavor decay electrons in $|y| < 0.35$ in high-multiplicity $d + \text{Au}$ events with respect to a combined data and theory $pp$ Ref. [15]. Recently, PHENIX also measured a significant enhancement of heavy-flavor production via single-muon detection at backward rapidity (the $\text{Au}$-going direction), and a suppression at forward rapidities (the $d$-going direction) [16]. This measured difference in heavy-flavor production between forward and backward rapidities is significantly larger than predicted by leading-order perturbative QCD calculations with nuclear PDFs [17]. In $p + \text{Pb}$ collisions at the LHC, the ALICE Collaboration measured the production of the $D$ meson in the $-0.96 < y < 0.04$ interval and found it to be, within uncertainties, compatible with $pp$ data scaled by the number of binary $NN$ collisions, over a large transverse momentum ($p_T$) range [18]. The LHC results are well described by theoretical calculations that do not require a deconfined medium to be formed in the collision. This supports the idea that the $D$ meson suppression at high $p_T$ observed in $\text{Pb} + \text{Pb}$ collisions by the ALICE Collaboration [19] is due to parton interactions with the deconfined medium. While measurements, both at RHIC and LHC, support that most of the suppression observed in $AA$ collisions is due to partonic energy loss, the details of the phenomena affecting open charm in $p + A$ and $AA$ collisions are still to be understood.

The production of $B$ mesons was studied at the LHC in proton-proton ($pp$) collisions at $\sqrt{s} = 7$ TeV over wide $p_T$ and rapidity intervals by CMS [20–22], ATLAS [23], and LHCb [24]. In $\text{Pb} + \text{Pb}$ collisions, CMS measured the non-prompt $J/\psi$ from $B$ hadron decays at $\sqrt{s_{NN}} = 2.76$ TeV.
[25] and observed a strong suppression with respect to the hypothesis of binary scaling. In this Letter, we extend the study of heavy-quark production in $p(d)+A$ collisions by performing the first measurement of exclusive $B$ meson decays in $p+Pb$ collisions.

The $B$ mesons are measured in a region $|y_{\text{lab}}|<2.4$ via the full reconstruction of their decay channels: $B^+ \rightarrow J/\psi K^+ \rightarrow \mu^+ \mu^- K^+$ with branching fraction $B = (6.12 \pm 0.19) \times 10^{-5}$, $B^0 \rightarrow J/\psi K^+\pi^- \rightarrow \mu^+ \mu^- K^+\pi^-$ with $B = (5.24 \pm 0.24) \times 10^{-5}$, and $B^0_{s} \rightarrow J/\psi K^+ \rightarrow \mu^+ \mu^- K^+K^-$ with $B = (3.12 \pm 0.27) \times 10^{-5}$ [26]. As this analysis does not separate $B^+$ from $B^0$, $B^0_{s}$ from $B^0_{s}$, mesons are referred to generically as $B^+$, $B^0$, and $B^0_{s}$, respectively, for the purposes of reconstruction. For the final cross section values, the combined results are divided by 2 to obtain an average.

The CMS detector has excellent capabilities to reconstruct $B$ meson decays due to the highly efficient muon detection system and the high-resolution silicon tracker [27]. The data sample used in this analysis corresponds to an integrated luminosity of $(34.6 \pm 1.2) \text{ nb}^{-1}$ [28]. The direction of the proton beam was initially opposite to the positive direction of the CMS longitudinal axis [27], and it was reversed after 60% of the data were taken. The beam energies were 4 TeV for protons and 1.58 TeV per nucleon for lead nuclei, resulting in a nucleon-nucleon center-of-mass energy of $\sqrt{s_{NN}} = 5.02$ TeV. Because of the energy difference of the colliding beams, the nucleon-nucleon center-of-mass frame in $p+Pb$ collisions was not at rest with respect to the laboratory frame. The results presented here use the convention that the proton-going side corresponds to positive pseudorapidity. This implies that massless particles emitted at pseudorapidity $\eta_{\text{CM}}$ in the $NN$ center-of-mass frame are detected at $\eta_{\text{lab}} = \eta_{\text{CM}} + 0.465$.

A detailed description of the CMS experiment and coordinate system can be found in Ref. [27]. Only the detector subsystems most relevant for this analysis are described here. Charged particles (tracks) are reconstructed within the range $|\eta_{\text{lab}}|<2.5$ by using the silicon tracker detector, located in the 3.8 $T$ magnetic field of a superconducting solenoid. Muons are identified in the interval $|\eta_{\text{lab}}|<2.4$ with gas-ionization detectors made of three technologies: drift tubes, cathode strip chambers, and resistive plate chambers, embedded in the steel flux-return yoke of the magnet. The CMS apparatus also has extensive forward calorimetry, including two steel and quartz-fiber Cherenkov hadron forward (HF) calorimeters, which cover the range $2.9<|\eta_{\text{lab}}|<5.2$.

Events used in the measurement are collected with a trigger requiring the presence of a muon with $p_T>3\text{ GeV}/c$. To select inelastic hadronic interactions, the off-line analysis requires a coincidence of at least one of the HF calorimeter towers (with more than 3 GeV of total energy) from each side of the interaction point. Events are further required to have at least one reconstructed primary vertex, formed by at least two tracks, with a distance from the center of the nominal interaction region of less than 15 cm along the beam axis.

Several Monte Carlo (MC) simulated event samples are used to evaluate background components and signal efficiencies: specifically, (i) an inclusive (prompt and nonprompt) $J/\psi$ sample; (ii) a sample containing all $B$ mesons decaying into a $J/\psi$; (iii) a signal-only sample with the $B^+$, $B^0$, and $B^0_{s}$ decays included in the present analysis. First, proton-proton collisions are simulated with PYTHIA 6.424 [29] tune ZZ [30] and propagated through the CMS detector using the GEANT4 package [31]. The $B$ meson decays are simulated with the EVTGEN package [32], and final state photon radiation in the $B$ decays is simulated by PHOTOS [33]. Then, the PYTHIA events are embedded into simulated $p+Pb$ events produced by the HIJING generator version 1.383 [34], which is tuned to reproduce global event properties such as charged-hadron $p_T$ spectra and particle multiplicity.

Muons are required to be within the following kinematic region: $p_T^\mu > 3.3\text{ GeV}/c$ for $|\eta_{\text{lab}}| < 1.3$, total momentum $p^\mu > 2.9\text{ GeV}/c$ for $1.3 < |\eta_{\text{lab}}| < 2.2$, or $p_T^\mu > 1.5\text{ GeV}/c$ for $2.2 < |\eta_{\text{lab}}| < 2.4$ [35]. This acceptance selection is chosen so as to guarantee a single-muon detection probability exceeding about 10%. Two muons of opposite charge with an invariant mass within 150 $\text{MeV}/c^2$ of the world-average $J/\psi$ mass [26] are selected to reconstruct a $J/\psi$ candidate, with a mass resolution of typically $18-55\text{ MeV}/c^2$, degrading as a function of the dimuon rapidity. The $B$ meson candidates are formed by combining $J/\psi$ candidates with charged tracks. Without using particle identification, assumptions need to be made about the masses of the charged tracks. In calculating the mass of the $B^+$ candidates, the single charged particle is always assumed to have the mass of a kaon. In the $B^0$ case, two invariant-mass values are computed, corresponding to the two possible assignments of the kaon and pion masses to the two-track system. For $B^0$ candidates, the two charged tracks are always assumed to be kaons. Single track low $p_T$ thresholds of 0.9, 0.7, and 0.4 $\text{GeV}/c$ are applied in the $B^+$, $B^0$, and $B^0_{s}$ analyses, respectively, to reduce the combinatorial background, which is further minimized by additional selection criteria. In particular, $B$ candidates are selected according to the $\chi^2$ probability of the decay vertex (the probability for the $J/\psi$ muon tracks and the other charged track to point to a common vertex), the 3D flight distance (normalized by its uncertainty) between the primary and decay vertices, and the pointing angle, which is defined as the angle between the line connecting the primary and decay vertices and the momentum vector of the $B$ meson in the plane transverse to the beam direction. The selection is optimized for each meson species using a multivariate technique that uses the genetics algorithm [36], in order to maximize the statistical significance of the $B$ meson signals. In the $B^0$ and $B^0_{s}$ analyses, the invariant masses...
of the $K^+\pi^-$ and the $K^+K^-$ are required to be compatible with the masses of the $K^{0}\pi(892)$, $K^*(892)$ and the $\phi$ resonances, respectively. If more than one candidate in a given event survives all of the aforementioned selection criteria, the candidate with the best vertex $\chi^2$ probability is selected.

The raw yields of $B^+$, $B^0$, and $B^{0*}$ are extracted using a binned maximum likelihood fit to the $B$ meson invariant-mass distributions in the mass range $5 < m_B < 6$ GeV/$c^2$. The invariant-mass distributions of $B^+$, $B^0$, and $B^{0*}$ candidates in the $p_T$ regions $10–15$, $10–15$, and $10–60$ GeV/$c$, respectively, are shown in Fig. 1. In the case of $B^+$ and $B^0$, this choice corresponds to the lowest $p_T$ interval used in the analysis, while for $B^{0*}$ it is the only interval. The signal shape is modeled by two Gaussians with the same mean values (a free parameter in the fit) and different widths determined in MC simulations. The background is dominated by random combinations of prompt and nonprompt $J/\psi$ candidates with extra particles. This combinatorial background is modeled by a first-order polynomial in the $B^+$ and $B^0$ analyses, and by a second-order polynomial in the $B^{0*}$ analysis, as suggested by studies on the embedded inclusive $J/\psi$ sample. The background component shown as a crosshatched histogram and labeled as $B \rightarrow J/\psi X$ in Fig. 1 is due to misreconstructed $B$ meson decays that produce broad peak structures in the invariant-mass region below 5.4 GeV/$c^2$. As an example, in the $B^+$ analysis, a peaking background structure is created by $B^{0} \rightarrow J/\psi K^*(892)$ decays in which one decay product is lost in the $B$ candidate reconstruction. These background sources are studied with the embedded MC sample including all $B$ meson decays into final states with a $J/\psi$, and found to be well described by a superposition of four and two Gaussian functions in the $B^+$ and $B^0$ analyses, respectively. The resulting functional form, with the overall normalization left floating, is included in the global fit function. This background component is found to be negligible in the $B^{0*}$ analysis as a consequence of the selection on the mass of the $\phi$ candidate.

The $p_T$-differential production cross section of the various $B$ meson species is computed in each $p_T$ interval:

$$\frac{d\sigma}{dp_T}|_{|y_{lab}|<2.4} = \frac{1}{2}\Delta p_T (Acc\,\epsilon)_{|y_{lab}|<2.4} \mathcal{L}.$$ (1)

$N(p_T)_{|y_{lab}|<2.4}$ is the raw signal yield extracted in each $p_T$ interval of width $\Delta p_T$, $(Acc\,\epsilon)_{|y_{lab}|<2.4}$ represents the corresponding acceptance times efficiency, $\mathcal{L}$ is the integrated luminosity, and $B$ is the branching fraction of the decay chain. The factor 1/2 accounts for the fact that the yields were measured for particles and antiparticles added together, but the cross section is given for particles only. An analogous expression holds for the rapidity-differential cross section. The $(Acc\,\epsilon)$ correction factors are evaluated using PYTHIA+HIJING simulations in each $p_T$ and $|y_{lab}|$ interval, to account for the loss of signal due to the detector coverage, and to the trigger, reconstruction, and off-line selection. They vary, over the measured $p_T$ and $y$ intervals, from 9% to 37% (3% to 15%) for $B^+$ ($B^0$), and they equal 8% in the single bin used for the $B^{0*}$.

The cross sections are affected by several sources of systematic uncertainties arising from the signal extraction, acceptance and efficiency corrections, branching fractions, and integrated luminosity determination. The uncertainty from the fitting procedure (varying from 10% to 15% across all analysis intervals, for all three mesons) is evaluated by varying the probability distribution functions used to model the signal and background distributions. As an alternative combinatorial background shape, a second-order polynomial is used for $B^+$ and $B^0$, and a third-order polynomial for $B^{0*}$. The uncertainty on the signal is evaluated by considering three fit variations: (i) leaving free the width parameters, (ii) varying the width parameters by $\pm 20\%$ with respect to the MC value, and (iii) using only
The systematic uncertainties associated with the track reconstruction data and the embedded MC sample. The systematic uncertainty from a comparison of two-body and four-body decays is estimated as systematic uncertainty, and it is given in the case of $B^+$ mesons, $B^0$, $B_s^0$, and $B_c^+$, respectively [26].

In Fig. 2, the $p_T$-differential production cross sections of all three $B$ mesons measured in the interval $|y_{\text{lab}}| < 2.4$ are presented, with data points placed at the center of each bin. They are compared to the $p$ cross sections obtained from fixed-order plus next-to-leading-logarithm (FONLL) calculations [37], which reproduce the $B$ meson $p_T$-differential cross sections in $pp$ collisions at 7 TeV [20–24]. The individual cross sections are obtained by scaling the FONLL total beauty production [37–39] by the world-average production fractions of $B^+$, $B^0$, and $B_s^0$ (40.2%, 40.2%, and 10.5%, respectively [26]). The obtained $B^+$ FONLL reference is validated using published experimental cross sections measured in $pp$ collisions at $\sqrt{s} = 7$ TeV [20,23]. The FONLL predictions are scaled by $A(= 208)$, the atomic mass of the Pb nucleus, to account for the number of binary $NN$ collisions [42]. The FONLL uncertainties, which are larger than the experimental uncertainties, represent the quadratic sum of several variations made to the calculation: of the factorization and renormalization scales, of the $b$ quark mass, and of the uncertainty associated with PDFs (providing the largest contribution) [37–39]. The nuclear modification factor $R_{pA}$, shown in Fig. 3, is computed as

$$R_{pA} = \frac{d^3N_{pA}}{dp_t dy} / \frac{d^3N_{pp}}{dp_t dy}$$

FIG. 2. The $p_T$-differential production cross section of $B^+$ (left panel), $B^0$ (center panel), and $B_s^0$ (right panel) measured in $p + \text{Pb}$ collisions at $\sqrt{s_{NN}} = 5.02$ TeV. The vertical bars (boxes) correspond to statistical (systematic) uncertainties. The global systematic uncertainty, listed in each panel and not included in the data points, comprises the uncertainties in the integrated luminosity measurement and the $B$ meson branching fractions. Results are compared to FONLL calculations [37–39], scaled by the number of binary $NN$ collisions, represented by a continuous histogram. The dashed histograms represent the theoretical uncertainties for the FONLL reference.
FIG. 3. The nuclear modification factors $R_{p+\text{A}}^\text{FONLL}(p_T)$ of $B^+$ (left panel), $B^0$ (center panel), $B_s^0$ (right panel) measured in $p + \text{Pb}$ collisions at $\sqrt{s_{NN}} = 5.02$ TeV. The statistical and systematic uncertainties on the $p + \text{Pb}$ data are shown as bars and yellow boxes around the data points, respectively. The systematic uncertainties from the FONLL predictions are plotted separately as open blue boxes. The global systematic uncertainties are shown as full grey boxes at unity, and are not included in the data points.

$$R_{p+\text{A}}^\text{FONLL} = \frac{\left( \frac{d\sigma}{dp_T} \right)_{p+\text{Pb}}}{A \left( \frac{d\sigma}{dp_T} \right)_{\text{FONLL}}},$$

where the numerator is defined in Eq. (1) and the denominator is the corresponding theoretical calculation for $B$ meson production in $pp$ collisions at the same center-of-mass energy. The theoretical uncertainties represented by the open blue boxes in Fig. 3 are computed by recalculating $R_{p+\text{A}}^\text{FONLL}(p_T)$ with the upper and lower values of the FONLL predictions represented by dashed histograms in Fig. 2.

The nuclear modification factors of the three $B$ mesons do not show evidence for modification of $p + \text{Pb}$ data compared to the FONLL reference, in the considered $p_T$ range within the quoted uncertainties. No significant differences are observed between the three $B$ meson species. In the lowest $p_T$ interval measured, $R_{p+\text{A}}^\text{FONLL}(p_T)$ is $1.11 \pm 0.08(\text{stat}) \pm 0.17(\text{syst} + \text{Pb}) +0.33^{+0.25}_{-0.29}$ (syst FONLL), $0.99 \pm 0.15(\text{stat}) \pm 0.18(\text{syst} + \text{Pb}) +0.30^{+0.26}_{-0.26}$ (syst FONLL), and $0.93 \pm 0.18(\text{stat}) \pm 0.20(\text{syst} + \text{Pb}) +0.27^{+0.27}_{-0.24}$ (syst FONLL), for $B^+$, $B^0$, and $B_s^0$, respectively.

The production cross section of $B^+$ is also studied as a function of its rapidity in the center-of-mass frame ($y_{CM}$). The $y_{CM}$-differential cross section of $B^+$ in the interval

FIG. 4. (Left panel) The $y_{CM}$-differential production cross section of $B^+$ measured in $p + \text{Pb}$ collisions at $\sqrt{s_{NN}} = 5.02$ TeV. Vertical bars (the boxes) correspond to statistical (systematic) uncertainties. The listed global systematic uncertainty is not included in the data points. The result is compared to a FONLL calculation [37–39] represented by a continuous histogram. The dashed histograms represent the theoretical uncertainties for the FONLL reference. (Right panel) The nuclear modification factor $R_{p+\text{A}}^\text{FONLL}(y_{CM})$ of $B^+$ as a function of $y_{CM}$. The statistical and systematic uncertainties on the $p + \text{Pb}$ data are shown as bars and yellow boxes around the data points, respectively. The systematic uncertainty from the FONLL reference is plotted separately as open blue boxes. The global systematic uncertainty is shown as a full grey box at unity, and it is not included in the data points.
10 < p_T < 60 GeV/c is shown in Fig. 4 (left panel). In Fig. 4 (right panel), the rapidity dependence of the nuclear modification factor of B^+ is shown. No strong evidence of rapidity dependence of R_{FONLL}^{p,A} is observed within the uncertainties.

In summary, the first measurements of the B^+, B^0, and B^0_s meson production cross sections in p + Pb collisions at √s_{NN} = 5.02 TeV are presented. The mesons are measured in |y_{lab}| < 2.4 and 10 < p_T < 60 GeV/c via the reconstruction of one of their exclusive hadronic decay channels. Within the transverse momentum and rapidity ranges studied, no significant modifications are observed, considering the statistical and systematical uncertainties, when compared to pp FONLL calculations scaled by the number of incoherent nucleon-nucleon collisions. These results provide a baseline for the study of in-medium b quark energy loss in Pb + Pb collisions.

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<td>University of Wisconsin, Madison, USA</td>
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</tbody>
</table>
aDeceased.
bAlso at Vienna University of Technology, Vienna, Austria.
cAlso at CERN, European Organization for Nuclear Research, Geneva, Switzerland.
dAlso at State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China.
eAlso at Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France.
fAlso at National Institute of Chemical Physics and Biophysics, Tallinn, Estonia.
gAlso at Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia.
hAlso at Universidade Estadual de Campinas, Campinas, Brazil.
iAlso at Centre National de la Recherche Scientifique (CNRS) - IN2P3, Paris, France.
jAlso at Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France.
kAlso at Joint Institute for Nuclear Research, Dubna, Russia.
lAlso at Helwan University, Cairo, Egypt.
mAlso at Zewail City of Science and Technology, Zewail, Egypt.
Also at Ain Shams University, Cairo, Egypt.
Also at British University in Egypt, Cairo, Egypt.
Also at Beni-Suef University, Bani Sweif, Egypt.
Also at Université de Haute Alsace, Mulhouse, France.
Also at Tbilisi State University, Tbilisi, Georgia.
Also at University of Hamburg, Hamburg, Germany.
Also at Brandenburg University of Technology, Cottbus, Germany.
Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary.
Also at Eötvös Loránd University, Budapest, Hungary.
Also at University of Debrecen, Debrecen, Hungary.
Also at Wigner Research Centre for Physics, Budapest, Hungary.
Also at University of Visva-Bharati, Santiniketan, India.
Also at King Abdulaziz University, Jeddah, Saudi Arabia.
Also at University of Ruhuna, Matara, Sri Lanka.
Also at Isfahan University of Technology, Isfahan, Iran.
Also at University of Tehran, Department of Engineering Science, Tehran, Iran.
Also at Plasma Physics Research Center, Science and Research Branch, Islamic Azad University, Tehran, Iran.
Also at Università degli Studi di Siena, Siena, Italy.
Also at Purdue University, West Lafayette, USA.
Also at International Islamic University of Malaysia, Kuala Lumpur, Malaysia.
Also at Malaysian Nuclear Agency, MOSTI, Kajang, Malaysia.
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Also at St. Petersburg State Polytechnical University, St. Petersburg, Russia.
Also at National Research Nuclear University ‘Moscow Engineering Physics Institute’ (MEPhI), Moscow, Russia.
Also at Faculty of Physics, University of Belgrade, Belgrade, Serbia.
Also at Facoltà Ingegneria, Università di Roma, Roma, Italy.
Also at National Technical University of Athens, Athens, Greece.
Also at Scuola Normale e Sezione dell’INFN, Pisa, Italy.
Also at University of Athens, Athens, Greece.
Also at Warsaw University of Technology, Institute of Electronic Systems, Warsaw, Poland.
Also at Institute for Theoretical and Experimental Physics, Moscow, Russia.
Also at Albert Einstein Center for Fundamental Physics, Bern, Switzerland.
Also at Adiyaman University, Adiyaman, Turkey.
Also at Mersin University, Mersin, Turkey.
Also at Cag University, Mersin, Turkey.
Also at Piri Reis University, Istanbul, Turkey.
Also at Gaziosmanpasa University, Tokat, Turkey.
Also at Ozyegin University, Istanbul, Turkey.
Also at Izmir Institute of Technology, Izmir, Turkey.
Also at Mimar Sinan University, Istanbul, Istanbul, Turkey.
Also at Marmara University, Istanbul, Turkey.
Also at Kafkas University, Kars, Turkey.
Also at Yildiz Technical University, Istanbul, Turkey.
Also at Hacettepe University, Ankara, Turkey.
Also at Rutherford Appleton Laboratory, Didcot, United Kingdom.

Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom.

Also at Instituto de Astrofísica de Canarias, La Laguna, Spain.

Also at Utah Valley University, Orem, USA.

Also at University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia.

Also at Argonne National Laboratory, Argonne, USA.

Also at Erzincan University, Erzincan, Turkey.

Also at Texas A&M University at Qatar, Doha, Qatar.

Also at Kyungpook National University, Daegu, Korea.