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CLT for Circular beta-Ensembles at high temperature

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ABSTRACT

We consider the macroscopic large N limit of the Circular beta-Ensemble at high temperature, and its weighted version as well, in the regime where the inverse temperature scales as β/N for some parameter $\beta > 0$. More precisely, in the limit $N \rightarrow \infty$, the equilibrium measure of this particle system is described as the unique minimizer of a functional which interpolates between the relative entropy ($\beta = 0$) and the weighted logarithmic energy ($\beta = \infty$). The purpose of this work is to show that the fluctuation of the empirical measure around the equilibrium measure converges towards a Gaussian field whose covariance structure interpolates between the Lebesgue L^2 ($\beta = 0$) and the Sobolev $H^{1/2}$ ($\beta = \infty$) norms. We furthermore obtain a rate of convergence for the fluctuations in the W_2 metric. Our proof uses the normal approximation result of Lambert et al. [24], the Coulomb transport inequality of Chafaï et al. [8], and a spectral analysis for the operator associated with the limiting covariance structure.

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1. Introduction and statement of the results

Let $\mathbb{T} := \mathbb{R}/2\pi\mathbb{Z}$ be the one-dimensional torus that we equip with the metric $(x, y) \mapsto |e^{ix} - e^{iy}| = |2 \sin(\frac{x-y}{2})|$. Given an inverse temperature parameter $\beta > 0$, the Circular beta-Ensemble is a celebrated particle system from random matrix theory of N particles on \mathbb{T} with distribution

$$\frac{1}{Z_N} \prod_{i < j} |e^{ix_i} - e^{ix_j}|^\beta \prod_{i=1}^N \frac{dx_i}{2\pi}$$

where $Z_N > 0$ is a normalization constant. This corresponds to the eigenvalues distribution of a unitary Haar distributed random matrix when $\beta = 2$. The macroscopic behavior of this particle system as $N \rightarrow \infty$ is well-known: the empirical measure

$$\mu_N := \frac{1}{N} \sum_{i=1}^N \delta_{x_i} \tag{1}$$

converges almost surely (a.s.) weakly towards the uniform measure $\frac{dx}{2\pi}$ on \mathbb{T} . The fluctuations of the particle system around the uniform measure can be described as well: for any smooth enough test function $\psi : \mathbb{T} \rightarrow \mathbb{R}$ satisfying $\int_{\mathbb{T}} \psi \frac{dx}{2\pi} = 0$, Johansson [19] proved¹ the central limit theorem (CLT)

$$N \int \psi d\mu_N = \sum_{i=1}^N \psi(x_i) \xrightarrow[N \rightarrow \infty]{\text{law}} \mathcal{N}(0, \frac{2}{\beta} \|\psi\|_{\mathbb{H}^{1/2}}^2), \tag{2}$$

where the Sobolev semi-norm $\|\cdot\|_{\mathbb{H}^{1/2}}$ is defined by

¹ More precisely, the CLT in [19] is stated for $\beta = 2$, in which case it is equivalent to the strong Szegő theorem for Toeplitz determinants, see for example [35, Chapter 6] or [10] for comprehensive expositions of this celebrated result. However, it is straightforward to check that the method of [19] still applies for any fixed $\beta > 0$ provided that the test function ψ is $C^{1+\alpha}$ for some $\alpha > 0$. See also [22, Theorem 1.2] for a generalization to the mesoscopic scale. Let us also stress that, although one may believe this CLT holds true as soon as $\|\psi\|_{\mathbb{H}^{1/2}} < \infty$, a counterexample has been provided in [22] when $\beta = 4$.

$$\|\psi\|_{\mathbf{H}^{1/2}}^2 := 2 \sum_{k=1}^{\infty} k |\hat{\psi}_k|^2.$$

Here and in what follows $\hat{\psi}_k := \int_{\mathbb{T}} \psi(x) e^{-ikx} \frac{dx}{2\pi}$ are the usual Fourier coefficients of ψ .

The aim of this work is to provide similar statements at high temperature, namely when β goes to zero as $N \rightarrow \infty$. Notice first that if we take $\beta = 0$, which corresponds to the infinite temperature setting, then the x_i 's are independent random variables uniformly distributed on \mathbb{T} . Thus the law of large numbers yields the a.s. weak convergence $\mu_N \rightarrow \frac{dx}{2\pi}$ as $N \rightarrow \infty$ and the classical CLT states that, for any L^2 function $\psi : \mathbb{T} \rightarrow \mathbb{R}$ satisfying $\hat{\psi}_0 = 0$,

$$\sqrt{N} \int \psi d\mu_N = \frac{1}{\sqrt{N}} \sum_{i=1}^N \psi(x_i) \xrightarrow[N \rightarrow \infty]{\text{law}} \mathcal{N}(0, \|\psi\|_{L^2}^2) \tag{3}$$

where the L^2 norm reads

$$\|\psi\|_{L^2}^2 := \int_{\mathbb{T}} \psi(x) \frac{dx}{2\pi} = 2 \sum_{k=1}^{\infty} |\hat{\psi}_k|^2.$$

Notice the difference of normalization between (2) and (3).

As we shall see, there is a critical temperature regime of temperature where the variance structure of the fluctuations interpolates between the Lebesgue L^2 and Sobolev $\mathbf{H}^{1/2}$ (semi-)norms, and this happens when β is of order $1/N$. Thus, from now we consider the particle system where we rescale the inverse temperature parameter as $\beta \mapsto 2\beta/N$, the factor 2 being cosmetic. We also consider the case where the particle system is subject to an external potential V and will show that the limiting variance depends on V in a non-trivial way. In contrast, in the usual fixed temperature setting, the variance is expected to depend only on the support of the equilibrium measure.

The study of random matrix ensembles at *high temperature* (i.e. with an interaction strength of order $1/N$) was initiated by Allez et al. [2] who described explicitly the crossover for the density of state from the Wigner semicircle law to the Gaussian law. There are also several results about eigenvalues fluctuations in this regime [4,27,29,28,23] whose study is motivated by the transition from random matrix to Poisson statistics, which is considered to be instrumental to describe the Anderson localization phenomenon. In particular Trinh [37] and Nakano and Trinh [27] obtained a CLT for the linear statistics of the Gaussian-beta-ensembles at this temperature regime, relying on the Dumitriu and Edelman [14] tridiagonal matrix representation for this particle system, although the limiting variance is not explicit. The asymptotic behavior of the largest eigenvalue of the Gaussian beta-ensembles at high temperature has been recently investigated in [30]. Moreover, in [36], the asymptotic behavior of the generalized free

energy of the Toda chain has also been related with certain statistics of the Dumitriu–Edelman model in the high temperature regime. There are also a few results available in higher dimension for Coulomb gases [33,1] in this regime. Here we chose to focus instead on beta-ensembles on \mathbb{T} ; that \mathbb{T} is compact yields several technical simplifications in the proofs and a simple formula for the limiting variance. However, let us mention that one could adapt our approach to tackle the setting of the Gaussian-beta-ensembles, and the beta-ensembles on \mathbb{R} with a general potential as well, and provide an explicit formula for the limiting variance similar to the one that we will derive below.

Let us also mention that a similar regime has been partially studied by Guionnet and Bodineau [17] for a two-component plasma model confined in a 2D box.

We now present the particle system we investigate and our main results.

The particle system of interest. For any $\beta > 0$ and any continuous potential $V : \mathbb{T} \rightarrow \mathbb{R}$, we consider N random interacting particles on \mathbb{T} with joint probability distribution

$$d\mathbb{P}_N(x_1, \dots, x_N) := \frac{1}{Z_N} \prod_{i < j} |e^{ix_i} - e^{ix_j}|^{\frac{2\beta}{N}} \prod_{i=1}^N e^{-V(x_i)} \frac{dx_i}{2\pi} \quad (4)$$

where $Z_N > 0$ is a normalization constant (which depends on the parameters $\beta > 0$ and V). In the following we set

$$\mu_0^V(dx) := e^{-V(x)} \frac{dx}{2\pi} \quad (5)$$

and, without loss of generality (by adding a constant to V if necessary), we assume that μ_0^V is a probability measure on \mathbb{T} . If we introduce the discrete logarithmic energy of a configuration $\mathbf{x} = (x_1, \dots, x_N) \in \mathbb{T}^N$,

$$\mathcal{H}(\mathbf{x}) := \sum_{1 \leq i < j \leq N} \log |e^{ix_i} - e^{ix_j}|^{-1}$$

then (4) takes the form

$$d\mathbb{P}_N(\mathbf{x}) = \frac{1}{Z_N} \exp \left\{ -\frac{2\beta}{N} \mathcal{H}(\mathbf{x}) \right\} (\mu_0^V)^{\otimes N}(d\mathbf{x}),$$

which is the Gibbs measure associated with the energy interaction \mathcal{H} at inverse temperature $2\beta/N$ with reference measure $(\mu_0^V)^{\otimes N}$. This particle system has a physical interpretation: we can observe that $\mathcal{H}(\mathbf{x}) = \sum_{i < j} g(x_i - x_j)$ where g can be written as the restriction $g(x) = G(x, 0)$ of the Green function G of the two-dimensional torus $\mathbb{T} \times \mathbb{T}$, that is $\Delta G = -2\pi(\delta_0 - 1)$ on $\mathbb{T} \times \mathbb{T}$ in the distributional sense, see e.g. [6]. Thus \mathbb{P}_N describes a gas of N unit charges, interacting according to the laws of electrostatic on the two-dimensional torus but constrained to stay on $\mathbb{T} \simeq \mathbb{T} \times \{0\} \subset \mathbb{T} \times \mathbb{T}$, in

presence of an external potential V , at temperature $N/(2\beta)$. As we shall see below, one of the interesting reasons to consider the statistical properties of a Coulomb gas at high temperature is that there is a subtle competition between the energy and entropy of the gas which results in non-trivial global fluctuations. This fact is somewhat surprising knowing that for any $\beta \geq 0$, the local fluctuations of the Coulomb gas (4) are described by a homogeneous Poisson point process whose intensity is related to μ_β^V , [23].

Macroscopic behavior. First, we discuss the large N limit of the empirical measure μ_N , see (1), when the x_i 's are distributed according to \mathbb{P}_N . If μ lies in the space $\mathcal{M}_1(\mathbb{T})$ of probability measures on \mathbb{T} , define its logarithmic energy by

$$\begin{aligned} \mathcal{E}(\mu) &:= \iint \log \left| \sin \left(\frac{x-y}{2} \right) \right|^{-1} \mu(dx)\mu(dy) \in [0, +\infty] \\ &= \iint \log \frac{1}{|e^{ix} - e^{iy}|} \mu(dx)\mu(dy) + \log 2. \end{aligned} \tag{6}$$

Moreover, given any $\mu, \nu \in \mathcal{M}_1(\mathbb{T})$, the relative entropy of μ with respect to ν is given by

$$\mathcal{K}(\mu|\nu) := \int \log \left(\frac{d\mu}{d\nu} \right) d\mu \in [0, +\infty]$$

when μ is absolutely continuous with respect to ν ; set $\mathcal{K}(\mu|\nu) := +\infty$ otherwise. The functional of interest here is $F_\beta^V : \mathcal{M}_1(\mathbb{T}) \rightarrow [0, +\infty]$ defined by

$$F_\beta^V(\mu) := \beta \mathcal{E}(\mu) + \mathcal{K}(\mu|\mu_0^V). \tag{7}$$

Note that when $F_\beta^V(\mu)$ is finite, then μ is absolutely continuous and, if $\mu(dx) = \mu(x)dx$, then we can alternately write

$$F_\beta^V(\mu) = \beta \mathcal{E}(\mu) + \int V d\mu + \int \log \mu d\mu + \log(2\pi).$$

In particular, when μ has a density and $\int \log \mu d\mu < \infty$, we see that after scaling the potential with $\beta > 0$,

$$F_\infty^V(\mu) := \lim_{\beta \rightarrow \infty} \frac{1}{\beta} F_\beta^V(\mu) = \mathcal{E}(\mu) + \int V d\mu \tag{8}$$

is the celebrated weighted logarithmic energy from potential theory [34]. The next result can be extracted from the literature.

Theorem 1.1. *Let $\beta \geq 0$ and assume $V : \mathbb{T} \rightarrow \mathbb{R}$ is continuous.*

- (a) The functional $F_\beta^V : \mathcal{M}_1(\mathbb{T}) \rightarrow [0, +\infty]$ has compact level sets $\{F_\beta \leq \alpha\}$, $\alpha \in \mathbb{R}$, and is strictly convex. In particular it has a unique minimizer μ_β^V on $\mathcal{M}_1(\mathbb{T})$.
- (b) The sequence (μ_N) satisfies a large deviation principle in $\mathcal{M}_1(\mathbb{T})$ equipped with its weak topology at speed βN with rate function $\mu \mapsto F_\beta(\mu) - F_\beta(\mu_\beta^V)$. In particular,

$$\mu_N \xrightarrow[N \rightarrow \infty]{\text{a.s.}} \mu_\beta^V$$

in the probability space $\otimes_N(\mathbb{T}^N, \mathfrak{B}(\mathbb{T})^{\otimes N}, \mathbb{P}_N)$.

When $\beta = 0$, this is Sanov’s theorem for i.i.d. random variables and elementary properties of the relative entropy, see e.g. [11]. Moreover, the unique minimizer of F_0^V is given by (5) and hence the notation is consistent. In the case where $\beta > 0$, statement (a) is classical (see e.g. the proof of Proposition 2.1 below) and (b) can be found in [5,16]. In fact, statement (a) of the theorem is also true for weaker regularity assumptions on V and also when $\beta = \infty$. Moreover, if one considers back the fixed temperature setting by taking the particle system (4) after the scaling $\beta \mapsto N\beta$ and $V \mapsto NV$, then statement (b) holds true at the same speed with rate function $F_\infty^V - F_\infty(\mu_\infty^V)$, see [18,3].

We will derive several properties for μ_β^V in Section 2 but let us already mention that, due to the rotational invariance, the equilibrium measure μ_β^0 for $V = 0$ is the uniform probability measure $\frac{dx}{2\pi}$ on \mathbb{T} for every $\beta \in [0, \infty]$. For a general potential V , we shall see that μ_β^V has a bounded density that is larger than a positive constant and is essentially as smooth as V is.

Macroscopic fluctuations. Our main result is a central limit theorem (CLT) for the random signed measure

$$\nu_N := \sqrt{N}(\mu_N - \mu_\beta^V)$$

tested against sufficiently smooth functions, with an explicit upper bound on the rate of convergence in the Wasserstein W_2 metric; the latter is defined for random variables X, Y taking values in \mathbb{R}^d by

$$W_2(X, Y) := \inf_{Z \in \Pi(X, Y)} \sqrt{\mathbb{E}[\|Z_1 - Z_2\|^2]} \tag{9}$$

where the infimum is taken over all random variables $Z = (Z_1, Z_2)$ with $Z_1 \stackrel{\text{law}}{=} X$ and $Z_2 \stackrel{\text{law}}{=} Y$.

To state the result, let us also write μ_β^V for the density of the equilibrium measure, so that $d\mu_\beta^V(x) = \mu_\beta^V(x)dx$, and introduce the operator \mathcal{L} defined by

$$-\mathcal{L}\phi = \phi'' + 2\pi\beta H(\mu_\beta^V \phi') + (\log \mu_\beta^V)' \phi' \tag{10}$$

which acts formally on the space $L^2(\mathbb{T})$ of real-valued square integrable functions on \mathbb{T} equipped with the scalar product

$$\langle f, g \rangle_{L^2} := \int_{\mathbb{T}} f(x)g(x) \frac{dx}{2\pi}.$$

Here H stands for the Hilbert transform defined on $L^2(\mathbb{T})$ by

$$H\psi(x) := -\text{p.v.} \int_{\mathbb{T}} \frac{\psi(t)}{\tan\left(\frac{x-t}{2}\right)} \frac{dt}{2\pi} \tag{11}$$

where p.v. is the Cauchy principal value, that is the limit as $\varepsilon \rightarrow 0$ of this integral restricted to the integration domain $|e^{ix} - e^{it}| > \varepsilon$. Note that when $\beta = 0$ the operator \mathcal{L} corresponds to the Sturm-Liouville operator $\mathcal{L}\phi = -\phi'' + V'\phi'$. As we shall see from Proposition 4.3 below, for any $\beta > 0$ the operator \mathcal{L} is well-defined and positive on the Sobolev-type space

$$\mathbb{H} := \left\{ \psi \in L^2(\mathbb{T}) : \psi' \in L^2(\mathbb{T}), \int \psi d\mu_{\beta}^V = 0 \right\}, \tag{12}$$

which is an Hilbert space once equipped with the inner-product

$$\langle \phi, \psi \rangle_{\mathbb{H}} := \int \phi' \psi' d\mu_{\beta}^V,$$

and moreover that its inverse \mathcal{L}^{-1} is trace-class on \mathbb{H} .

The central result of this work is that ν_N converges, in the sense of finite dimensional distributions, to a Gaussian process on \mathbb{H} with covariance operator \mathcal{L}^{-1} .

Theorem 1.2 (CLT). *Let $\beta > 0$ and $V \in \mathcal{C}^{3,1}(\mathbb{T})$. Assume $\psi \in \mathcal{C}^{2\gamma+1}(\mathbb{T})$ for some integer $\gamma \geq 2$ and that $\int \psi d\mu_{\beta}^V = 0$. Then we have*

$$\nu_N(\psi) = \frac{1}{\sqrt{N}} \sum_{i=1}^N \psi(x_i) \xrightarrow[N \rightarrow \infty]{\text{law}} \mathcal{N}(0, \sigma_{\beta}^V(\psi)^2)$$

where the variance is given by

$$\sigma_{\beta}^V(\psi)^2 := \langle \psi, \mathcal{L}^{-1}\psi \rangle_{\mathbb{H}} = \int \psi' (\mathcal{L}^{-1}\psi)' d\mu_{\beta}^V. \tag{13}$$

Moreover, there exists $C = C(\beta, V, \psi) > 0$ such that

$$W_2\left(\nu_N(\psi), \mathcal{N}(0, \sigma_{\beta}^V(\psi)^2)\right) \leq C \sqrt{\frac{\log N}{N^{\frac{\gamma-1}{\gamma+1}}}}.$$

Of course the theorem still holds for a general $\psi \in \mathcal{C}^{2\gamma+1}(\mathbb{T})$ after replacing ψ by $\psi - \int \psi \, d\mu_\beta^V$, including in the limiting variance (13).² When $V = 0$, we can obtain an explicit formula for the limiting variance.

Lemma 1.3. *When $V = 0$, we have*

$$\sigma_\beta^0(\psi)^2 = 2 \sum_{k=1}^\infty \frac{1}{1 + \beta/k} |\hat{\psi}_k|^2.$$

This identity follows from the fact that, using the invariance by rotation, it is easy to diagonalize the operator \mathcal{L} – see the identity (52) below. Indeed, in this setting we have $-\mathcal{L}\phi = \phi'' + \beta H(\phi')$ and the eigenfunctions are given by the Fourier basis $\phi_j(x) = e^{ijx}$ since $\mathcal{L}\phi_j = (j^2 + \beta|j|)\phi_j$ for every $j \in \mathbb{Z}$.

Recalling (2) and (3), observe that $\sigma_\beta^0(\psi)^2 \rightarrow \|\psi\|_{L^2}^2$ as $\beta \rightarrow 0$ and that $\beta\sigma_\beta^0(\psi)^2 \rightarrow \|\psi\|_{\mathbb{H}^{1/2}}^2$ as $\beta \rightarrow \infty$; the factor 2 disappears due to the change of scale we made for temperature. In this sense $\sigma_\beta^0(\psi)$ interpolates between the Lebesgue L^2 and the Sobolev $\mathbb{H}^{1/2}$ (semi-)norm. In Section 8, we establish that for a general potential, we also have $\sigma_\beta^V(\psi)^2 \rightarrow \|\psi\|_{L^2(\mu_0^V)}^2$ as $\beta \rightarrow 0$ (see Proposition 8.2). We will also provide a sufficient condition on the equilibrium measure μ_β^V so that $\beta\sigma_\beta^V(\psi)^2 \rightarrow \|\psi\|_{\mathbb{H}^{1/2}}^2$ as well as $\beta\sigma_\beta^{\beta V}(\psi)^2 \rightarrow \|\psi\|_{\mathbb{H}^{1/2}}^2$ as $\beta \rightarrow \infty$ (see Proposition 8.3). This establishes that the Gaussian process which appears in Theorem 1.2 interpolates from a white noise (Poisson statistics) to a $\mathbb{H}^{1/2}$ noise (random matrix statistics). This also shows that the fluctuations become universal, in the sense that they do not depend on V , only when $\beta = \infty$.

Remark 1.1. Let us observe that the rate of convergence in Theorem 1.2 does not depend on the smoothness of V , but it improves with the regularity of the test function. Moreover, if $\psi \in \mathcal{C}^\infty(\mathbb{T})$, we have

$$W_2\left(\nu_N(\psi), \mathcal{N}(0, \sigma_\beta^V(\psi)^2)\right) \leq C\sqrt{\frac{\log N}{N}}.$$

We expect this rate to be optimal, maybe up to the factor $\sqrt{\log N}$.

The proof of Theorem 1.2 is deferred to Section 4 and relies on a normal approximation technique introduced in [24], which is inspired from Stein’s method; see Theorem 4.5 below. In [24] this method has been used to investigate the rate of convergence of the fluctuations for beta-Ensembles on \mathbb{R} at fixed temperature. There is a substantial technical difference in the analysis which arises in the high temperature regime due to the fact that the operator \mathcal{L} has an extra Sturm-Liouville component. In particular, the spectral

² Note the operator \mathcal{L}^{-1} is only defined on the Hilbert space \mathbb{H} , see (12).

properties of \mathcal{L} are quite different and this yields changes in the rate of convergence as well as in the limiting variance.

Stein’s method has also been used previously in the context of random matrix theory to investigate the rate of convergence for linear statistics of random matrices from the classical compact groups [15,12,13] and for the Circular beta-Ensemble at fixed temperature [38]. There are also results from Chatterjee [9] on linear statistics of Wigner matrices which are valid under strong assumptions on the law of the entries and from Johnson [20] on the eigenvalues of random regular graphs. For a comprehensive introduction to Stein’s method which includes several applications, we refer to the survey [32].

On the road to establish the CLT, we prove the following concentration inequality which may be of independent interest: let $W_1(\mu, \nu)$ be the Wasserstein-Kantorovich distance of order 1 between $\mu, \nu \in \mathcal{M}_1(\mathbb{T})$, defined by

$$W_1(\mu, \nu) := \inf_{\pi \in \Pi(\mu, \nu)} \iint |e^{ix} - e^{iy}| \pi(dx, dy) = \sup_{\|f\|_{\text{Lip}} \leq 1} \int f d(\mu - \nu) \tag{14}$$

where $\Pi(\mu, \nu)$ is the set of probability measures on $\mathbb{T} \times \mathbb{T}$ with respective marginals μ and ν ; the second identity is known as the Kantorovich-Rubinstein dual representation for W_1 , where the supremum is taken over Lipschitz functions $\mathbb{T} \rightarrow \mathbb{R}$ with Lipschitz constant at most one.

Theorem 1.4 (Concentration). *Let $\beta > 0$ and assume $V : \mathbb{T} \rightarrow \mathbb{R}$ has a weak derivative V' in $L^2(\mathbb{T})$. Then, there exists $C = C(\mu_\beta^V) > 0$ such that, for every $N \geq 10$ and $r > 0$,*

$$\mathbb{P}_N \left(W_1(\mu_N, \mu_\beta^V) > r \right) \leq e^{-\beta(\frac{1}{8\pi}Nr^2 - 5 \log N - C)}.$$

We have an explicit expression for the constant C in terms of μ_β^V in (30). In particular, when $V = 0$, this upper bound holds with $C = 2 \log 2 + 17/2 + \pi^{-1} \simeq 10.2$, which does not depend on β .

In particular, this yields together with Borel-Cantelli lemma that $W_1(\mu_N, \mu_\beta^V) \rightarrow 0$ a.s. for fixed $\beta > 0$ and, when $V = 0$, that $W_1(\mu_N, \frac{dx}{2\pi}) \rightarrow 0$ a.s. when β may depend on N as long as $\beta \gg N^{-1}$. For lower order temperature scales this should still be true but one needs to prove it differently; note also there is an interesting change of behavior for the partition function of the Gaussian-beta-ensemble around $\beta \sim N^{-1}$ pointed out in [29, Lemma 1.3].

The proof of the theorem follows the same strategy as the one of [8] and relies on their Coulomb transport inequality. Differences however arise due to the presence of the relative entropy in F_β^V . In particular, one needs to study the regularity of the potential of the equilibrium measure.

Organization of the paper. In section 2 we obtain preliminary results on the equilibrium measure μ_β^V and its logarithmic potential. Section 3 is devoted to the proof of

Theorem 1.4. In section 4, we provide the core of the proof of Theorem 1.2. In section 5, we obtain concentration estimates for error terms by means of Theorem 1.4. In section 6, we investigate spectral properties of the operator \mathcal{L} ; in particular we show that \mathcal{L}^{-1} is trace-class. In section 7, we study the regularity of the eigenfunctions of the operator \mathcal{L} so as to complete the proof of the main theorem. Finally, in Section 8, we investigate the behavior of the variance σ_β^V as $\beta \rightarrow 0$ (Poisson regime) as well as $\beta \rightarrow \infty$ (random matrix regime).

Notations, basic properties and conventions. From now, $\beta > 0$ is fixed. In the following, if η is a measure on \mathbb{T} , we will denote by $\eta(x)$ its density with respect to the Lebesgue measure dx when it exists. If $S \subset \mathbb{T}$ is a Borel set, we denote by $|S|$ its Lebesgue measure.

Recall that \mathbb{T} is equipped with the metric $(x, y) \mapsto |e^{ix} - e^{iy}|$ and denote for any $k \in \mathbb{N} := \{0, 1, 2, \dots\}$ and $0 < \alpha \leq 1$ by $\mathcal{C}^{k,\alpha}(\mathbb{T})$ the space of k -times differentiable functions on \mathbb{T} whose k -th derivative is α -Hölder continuous, or Lipschitz continuous when $\alpha = 1$. When $0 < \alpha < 1$ we also write \mathcal{C}^α instead of $\mathcal{C}^{0,\alpha}$, since there is not ambiguity, and put

$$\|\psi\|_{\mathcal{C}^\alpha} := \sup_{\substack{x,y \in \mathbb{T} \\ x \neq y}} \frac{|\psi(x) - \psi(y)|}{|e^{ix} - e^{iy}|^\alpha}, \quad \|\psi\|_{\text{Lip}} := \sup_{\substack{x,y \in \mathbb{T} \\ x \neq y}} \frac{|\psi(x) - \psi(y)|}{|e^{ix} - e^{iy}|}.$$

Note that, for any $0 < \alpha < 1$, we have $\|\psi\|_{\mathcal{C}^\alpha(\mathbb{T})} \leq 2\|\psi\|_{\text{Lip}}$.

We sometimes use as well the chordal metric

$$d_{\mathbb{T}}(x, y) := \inf_{k \in \mathbb{Z}} |x - y + 2k\pi|$$

instead of the reference metric since they are equivalent: $\frac{2}{\pi}d_{\mathbb{T}}(x, y) \leq |e^{ix} - e^{iy}| \leq d_{\mathbb{T}}(x, y)$ for any $x, y \in \mathbb{R}$. Moreover, since Rademacher’s theorem states that the Lipschitz constant for the metric $d_{\mathbb{T}}$ reads $\|f'\|_{L^\infty}$, we have

$$\|f'\|_{L^\infty} \leq \|f\|_{\text{Lip}} \leq \frac{\pi}{2}\|f'\|_{L^\infty}. \tag{15}$$

Recall that $\hat{\psi}_k = \int_{\mathbb{T}} \psi(x) e^{-ikx} \frac{dx}{2\pi}$ denotes the Fourier coefficient of $\psi \in L^1(\mathbb{T})$. Let $L^2_0(\mathbb{T}) = \{\psi \in L^2(\mathbb{T}) : \hat{\psi}_0 = 0\}$ and $H^m(\mathbb{T})$ be the Sobolev subspace of $L^2(\mathbb{T})$ of functions having their m -th first distributional derivatives in $L^2(\mathbb{T})$. We will also use at several instances the continuous embedding $H^{m+1}(\mathbb{T}) \subset \mathcal{C}^{m,1/2}(\mathbb{T})$ for $m \in \mathbb{N}$, sometimes known as the Sobolev-Hölder embedding theorem.

Note that as \mathbb{T} is compact, the Wasserstein-Kantorovich distances W_p are equivalent for all $p \geq 1$. The case W_1 is practical due to its dual-representation involving Lipschitz test functions (14), while W_2 is useful because its relate directly to the L^2 -distance (9).

Finally, we use the letter C for a positive constant which may vary from line to line, and which may depend only on $\beta > 0$ and on the potential V unless stated otherwise.

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2. Properties of the equilibrium measure

In this section we study the minimizer μ_β^V of F_β , see (7), and collect useful properties for later. Given $\mu \in \mathcal{M}_1(\mathbb{T})$, its logarithmic potential $U^\mu : \mathbb{T} \rightarrow [0, +\infty]$ is defined by

$$U^\mu(x) := \int \log \left| \sin \left(\frac{x-y}{2} \right) \right|^{-1} \mu(dy).$$

Proposition 2.1. *If $V : \mathbb{T} \rightarrow \mathbb{R}$ is a measurable and bounded function, then for any $\beta \geq 0$,*

- (a) F_β^V has a unique minimizer μ_β^V on $\mathcal{M}_1(\mathbb{T})$.
- (b) μ_β^V is absolutely continuous and there exists a $0 < \delta < 1$ such that

$$\delta \leq \frac{\mu_\beta^V(x)}{2\pi} \leq \delta^{-1} \quad a.e.$$

In particular, there exists $0 < \ell < 1$ such that $\ell \leq U^{\mu_\beta^V} \leq \ell^{-1}$ on \mathbb{T} .

- (c) *There exists a constant $C_\beta^V \in \mathbb{R}$ such that*

$$2\beta U^{\mu_\beta^V}(x) + V(x) + \log \mu_\beta^V(x) = C_\beta^V \quad a.e. \tag{16}$$

Part (c) of the proposition is usually referred as the Euler-Lagrange equation.

Remark 2.1. If $V = 0$, then μ_β^V is the uniform measure $\frac{dx}{2\pi}$ because of the rotational invariance. One can also check it satisfies (16) since, for any $x \in \mathbb{T}$,

$$U^{\frac{dx}{2\pi}}(x) = \int_{\mathbb{T}} \log \left| \frac{1 - e^{i(x-y)}}{2} \right|^{-1} \frac{dy}{2\pi} = \log 2. \tag{17}$$

Thus, the Euler-Lagrange constant reads $C_\beta^0 = 2\beta \log 2 - \log(2\pi)$.

Remark 2.2. Part (a) of the proposition follows from well known results. Although part (b) and (c) seem to be part of the folklore, we were not able to locate (b) and (c) proven in full details in the literature; the little subtlety is to take care of the sets where the density of μ_β^V may a priori vanish or be arbitrary close to zero due to the term $\log \mu_\beta^V$.

Proof of Proposition 2.1. It is known that both mappings $\mu \mapsto \mathcal{E}(\mu)$ and $\mu \mapsto \mathcal{K}(\mu|\mu_0^V)$ have compact level sets on $\mathcal{M}_1(\mathbb{T})$ and are strictly convex there, see [34,11], from which (a) directly follows. Moreover, since $F_\beta(\frac{dx}{2\pi}) < \infty$ we have $\mathcal{E}(\mu_\beta^V) < \infty$ and $\mathcal{K}(\mu_\beta^V|\mu_0^V) < \infty$. This implies μ_β^V is absolutely continuous because it has finite entropy with respect to an absolutely continuous probability measure.

Let $\mu_\beta^V : \mathbb{T} \rightarrow \mathbb{R}$ be any measurable function such that $\mu_\beta^V(dx) = \mu_\beta^V(x)dx$. We first claim that the Borel set $A_0 := \{x \in \mathbb{T} : \mu_\beta^V(x) = 0\}$ has null Lebesgue measure. Indeed, otherwise we could define $\eta := |A_0|^{-1}\mathbf{1}_{A_0}(x)dx \in \mathcal{M}_1(\mathbb{T})$ and obtain, for any $0 < \varepsilon < 1$,

$$\begin{aligned} F_\beta^V((1 - \varepsilon)\mu_\beta^V + \varepsilon\eta) &= F_\beta^V(\mu_\beta^V) \\ &+ \varepsilon \left(\int (2\beta U\mu_\beta^V + V)d(\eta - \mu) + \int \log \eta d\eta - \int \log \mu_\beta^V d\mu_\beta^V \right) \\ &+ \varepsilon \log \varepsilon + (1 - \varepsilon) \log(1 - \varepsilon) + \varepsilon^2 \beta \mathcal{E}(\mu_\beta^V - \eta). \end{aligned}$$

This yields in turn

$$F_\beta^V((1 - \varepsilon)\mu_\beta^V + \varepsilon\eta) = F_\beta^V(\mu_\beta^V) + \varepsilon(C + \log \varepsilon) + \mathcal{O}(\varepsilon^2)$$

when $\varepsilon \rightarrow 0$ for some $C \in \mathbb{R}$ and, since $\varepsilon(C + \log \varepsilon) + \mathcal{O}(\varepsilon^2)$ is negative for every $\varepsilon > 0$ small enough, this contradicts the fact that μ_β^V is the unique minimizer. Thus $|A_0| = 0$.

We next prove a weak form of (c). Let $\phi : \mathbb{T} \rightarrow \mathbb{R}$ be a measurable and bounded function satisfying $\int \phi d\mu_\beta^V = 0$. Then, for any real $|\varepsilon| \leq \|\phi\|_\infty^{-1}$, we have $(1 + \varepsilon\phi)\mu_\beta^V \in \mathcal{M}_1(\mathbb{T})$ and

$$\begin{aligned} F_\beta^V((1 + \varepsilon\phi)\mu_\beta^V) &= F_\beta^V(\mu_\beta^V) + \varepsilon \int (2\beta U\mu_\beta^V + V + \log \mu_\beta^V)\phi d\mu_\beta^V \\ &+ \varepsilon^2 \mathcal{E}(\phi \mu_\beta^V) + \int (1 + \varepsilon\phi) \log(1 + \varepsilon\phi) d\mu_\beta^V. \end{aligned}$$

By definition of μ_β^V , the mapping $\varepsilon \mapsto F_\beta^V((1 + \varepsilon\phi)\mu_\beta^V)$ has a unique minimum at $\varepsilon = 0$ and, since $\int (1 + \varepsilon\phi) \log(1 + \varepsilon\phi) d\mu_\beta^V = \mathcal{O}(\varepsilon^2)$, we obtain

$$\int (\beta U\mu_\beta^V + V + \log \mu_\beta^V)\phi d\mu_\beta^V = 0$$

for any such ϕ 's. If $\eta \in \mathcal{M}_1(\mathbb{T})$ has a bounded density ψ with respect to μ_β^V , then by taking $\phi := \psi - 1$ in the previous identity we obtain

$$\int (2\beta U\mu_\beta^V + V + \log \mu_\beta^V) d\eta = C_\beta^V := \int (2\beta U\mu_\beta^V + V + \log \mu_\beta^V) d\mu_\beta^V. \tag{18}$$

Now, if one assumes $A := \{x \in \mathbb{T} : 2\beta U\mu_\beta^V(x) + V(x) + \log \mu_\beta^V(x) > C_\beta^V\}$ has μ_β^V -positive measure, then by taking $\eta(dx) := \mu_\beta^V(A)^{-1}\mathbf{1}_A(x)\mu_\beta^V(dx)$ in (18) we reach a contradiction. Since the same holds after replacing $>$ by $<$ we obtain

$$2\beta U^{\mu_\beta^V} + V + \log \mu_\beta^V = C_\beta^V, \quad \mu_\beta^V\text{-a.e.} \tag{19}$$

We are now equipped to prove (b) and (c). Using that $U^{\mu_\beta^V} \geq 0$ on \mathbb{T} , we obtain from (19) that $\mu_\beta^V(x) \leq c e^{-V(x)}$ μ_β^V -a.e. for some $c > 0$, and thus the same holds true (Lebesgue)-a.e. In particular, since V is bounded by assumption, there exists $C > 0$ such that $\mu_\beta^V(x) \leq \frac{C}{2\pi}$ for a.e. $x \in \mathbb{T}$. This yields in turn with (17) that $U^{\mu_\beta^V}(x) \leq C U^{\frac{dx}{2\pi}}(x) = C \log 2$ on \mathbb{T} . Next, let $A_\kappa := \{x \in \mathbb{T} : \mu_\beta^V(x) \leq \kappa\}$ for any $0 < \kappa < 1$. If $\mu_\beta^V(A_\kappa) > 0$, then by taking the measure $\eta(dx) := (\mu_\beta^V(A_\kappa))^{-1} \mathbf{1}_{A_\kappa}(x) \mu_\beta^V(dx)$ in (18) we obtain

$$C_\beta^V \leq 2\beta C \log 2 + \|V\|_{L^\infty} + \log \kappa$$

and thus $\mu_\beta^V(A_\kappa) = 0$ for every $\kappa > 0$ small enough. Since we have already shown that $|A_0| = 0$, this means that $|A_\kappa| = 0$ for every $\kappa > 0$ small enough, and the first claim of (b) is proven. Since the function $x \mapsto \log |\sin(\frac{x}{2})|^{-1}$ is non-negative and integrable on \mathbb{T} , the second claim follows as well.

Finally, this yields that the equation (19) holds a.e. and thus (c) is proven. \square

Recall that one can extend the definition of the energy (6) to a class of signed measures. By [34, Lemma 1.8], $\mathcal{E}(\mu - \mu_\beta^V)$ makes sense and is non-negative for any $\mu \in \mathcal{M}_1(\mathbb{T})$ with $\mathcal{E}(\mu) < \infty$.

Corollary 2.2. *For any $\mu \in \mathcal{M}_1(\mathbb{T})$ satisfying $\mathcal{E}(\mu) < \infty$, we have*

$$F_\beta^V(\mu) - F_\beta^V(\mu_\beta^V) = \beta \mathcal{E}(\mu - \mu_\beta^V) + \mathcal{K}(\mu | \mu_\beta^V).$$

Proof. One can assume μ has a density which satisfies $\int \log \mu d\mu < \infty$ and $\mathcal{E}(\mu) < \infty$ since the identity is otherwise trivial. By integrating (16) against μ this yields

$$C_\beta^V = 2\beta \int U^{\mu_\beta^V} d\mu + \int V d\mu + \int \log \mu_\beta^V d\mu. \tag{20}$$

In particular, we obtain by taking $\mu = \mu_\beta^V$ and subtracting the resulting identity to (20),

$$\int V d(\mu - \mu_\beta^V) = 2\beta \mathcal{E}(\mu_\beta^V) - 2\beta \int U^{\mu_\beta^V} d\mu - \int \log \mu_\beta^V d\mu + \int \log \mu_\beta^V d\mu_\beta^V.$$

The latter identity plugged into $F_\beta^V(\mu) - F_\beta^V(\mu_\beta^V)$ yields the corollary. \square

We also describe the behavior as $\beta \rightarrow 0$ and $\beta \rightarrow \infty$ of the equilibrium measure.

Lemma 2.3. *If $V : \mathbb{T} \rightarrow \mathbb{R}$ is measurable and bounded, then we have the weak convergences*

$$\lim_{\beta \rightarrow 0} \mu_\beta^V = \mu_0^V \quad \text{and} \quad \lim_{\beta \rightarrow \infty} \mu_\beta^V = \mu_\infty^0 = \frac{dx}{2\pi}.$$

If we further assume V is lower semicontinuous and that μ_∞^V has a density which satisfies $\int \log \mu_\infty^V d\mu_\infty^V < \infty$, then we have the weak convergence

$$\lim_{\beta \rightarrow \infty} \mu_\beta^{\beta V} = \mu_\infty^V.$$

Note that V is lower semicontinuous and does not take the value $+\infty$ ensures that F_∞^V is lower semicontinuous and has a unique minimizer μ_∞^V on $\mathcal{M}_1(\mathbb{T})$, see [34].

Proof. First, since \mathcal{E} is positive, μ_β^V minimizes F_β , $\mathcal{K}(\mu_0^V | \mu_0^V) = 0$ and $\mathcal{E}(\mu_0^V) < \infty$, we have

$$\mathcal{K}(\mu_\beta^V | \mu_0^V) \leq F_\beta(\mu_\beta^V) \leq F_\beta(\mu_0^V) = \beta \mathcal{E}(\mu_0^V) \xrightarrow{\beta \rightarrow 0} 0.$$

Since $\mu \mapsto \mathcal{K}(\mu | \mu_0^V)$ has for unique minimizer μ_0^V and is lower semicontinuous on $\mathcal{M}_1(\mathbb{T})$, which is weakly compact, this implies the weak convergence $\mu_\beta^V \rightarrow \mu_0^V$ as $\beta \rightarrow 0$.

Next, recall that $\mu_\infty^0 = \frac{dx}{2\pi}$ is the unique minimizer of \mathcal{E} on $\mathcal{M}_1(\mathbb{T})$. Since

$$\beta \mathcal{E}(\mu_\beta^V) + \mathcal{K}(\mu_\beta^V | \mu_0^V) = F_\beta^V(\mu_\beta^V) \leq F_\beta^V\left(\frac{dx}{2\pi}\right) \leq \beta \mathcal{E}(\mu_\beta^V) + \mathcal{K}\left(\frac{dx}{2\pi} | \mu_0^V\right)$$

we obtain that $\mathcal{K}(\mu_\beta^V | \mu_0^V) \leq \mathcal{K}\left(\frac{dx}{2\pi} | \mu_0^V\right) = \int_{\mathbb{T}} V \frac{dx}{2\pi} < \infty$ for every $\beta > 0$, and moreover

$$\limsup_{\beta \rightarrow \infty} \mathcal{E}(\mu_\beta^V) = \limsup_{\beta \rightarrow \infty} \frac{1}{\beta} F_\beta^V(\mu_\beta^V) \leq \limsup_{\beta \rightarrow \infty} \frac{1}{\beta} F_\beta^V\left(\frac{dx}{2\pi}\right) = \mathcal{E}\left(\frac{dx}{2\pi}\right).$$

Since \mathcal{E} is lower semicontinuous on $\mathcal{M}_1(\mathbb{T})$, this similarly yields the weak convergence $\mu_\beta^V \rightarrow \frac{dx}{2\pi}$ as $\beta \rightarrow \infty$.

Finally, by observing that $F_\beta^{\beta V}(\mu) = \beta F_\infty^V(\mu) + \int \log(2\pi\mu) d\mu$ provided $\mu \in \mathcal{M}_1(\mathbb{T})$ has a density, we have

$$\begin{aligned} \beta F_\infty^V(\mu_\beta^{\beta V}) + \int \log(2\pi\mu_\beta^{\beta V}) d\mu_\beta^{\beta V} &= F_\beta^{\beta V}(\mu_\beta^{\beta V}) \\ &\leq F_\beta^{\beta V}(\mu_\infty^V) \\ &= \beta F_\infty^V(\mu_\infty^V) + \int \log(2\pi\mu_\infty^V) d\mu_\infty^V \end{aligned}$$

Thus, if $\int \log \mu_\infty^V d\mu_\infty^V < \infty$, after dividing by $\beta > 0$ and taking the limit as $\beta \rightarrow +\infty$, this implies

$$\limsup_{\beta \rightarrow \infty} F_\infty^V(\mu_\beta^{\beta V}) \leq F_\infty^V(\mu_\infty^V)$$

and the weak convergence $\mu_\beta^{\beta V} \rightarrow \mu_\infty^V$ as $\beta \rightarrow \infty$ is obtained as well. \square

Next, we study the regularity of the equilibrium measure and its potential. Recall the Hilbert transform H acting on the Hilbert space $L^2(\mathbb{T})$ is defined in (11). We can also define $H\mu$ for $\mu \in \mathcal{M}_1(\mathbb{T})$ as soon as it has a density $\mu(x)$. Note that H acts in a simple fashion on the Fourier basis: $H(1) = 0$ and, if $k \in \mathbb{N} \setminus \{0\}$,

$$H(e^{ikx}) = \frac{ie^{ikx}}{2\pi} \int_{\mathbb{T}} \frac{1 - e^{ik(t-x)}}{1 - e^{i(t-x)}} (1 + e^{i(t-x)}) dt = ie^{ikx}.$$

By taking the complex conjugate, this implies that for every $k \in \mathbb{Z}$,

$$H(e^{ikx}) = i \operatorname{sgn}(k)e^{ikx}$$

where we set $\operatorname{sgn}(0) := 0$. This yields that $H : L^2(\mathbb{T}) \rightarrow L^2_0(\mathbb{T})$ is a well-defined bounded operator with adjoint $H^* = -H$. Moreover, when restricted to $L^2_0(\mathbb{T})$, this turns H into an isometry which satisfies $H^{-1} = -H$. We will also use that this implies that for any $f \in H^1(\mathbb{T})$, Hf also belong to the Sobolev space $H^1(\mathbb{T})$ and that $(Hf)' = H(f')$. In the sequel, we will use these properties of the Hilbert transform at several instances.

Lemma 2.4. *If V is measurable and bounded, then $U^{\mu_\beta^V} \in H^1(\mathbb{T})$ and $(U^{\mu_\beta^V})' = \pi H\mu_\beta^V$.*

Proof. For any $\vartheta \in \mathcal{C}^1(\mathbb{T})$, by using the definition of the Cauchy principal value and doing an integration by part we obtain, for every $x \in \mathbb{T}$,

$$\begin{aligned} U^{\vartheta'}(x) &= \int_{\mathbb{T}} \log \left| \sin \left(\frac{x-t}{2} \right) \right|^{-1} \vartheta'(t) dt \\ &= -\text{p.v.} \int_{\mathbb{T}} \frac{1}{2 \tan \left(\frac{x-t}{2} \right)} \vartheta(t) dt \\ &= \pi H\vartheta(x). \end{aligned}$$

Next, using Fubini theorem and that H is a bounded operator on $L^2(\mathbb{T})$ satisfying $H^* = -H$, we obtain

$$\langle \vartheta', U^{\mu_\beta^V} \rangle_{L^2} = \langle U^{\vartheta'}, \mu_\beta^V \rangle_{L^2} = \langle \pi H\vartheta, \mu_\beta^V \rangle_{L^2} = -\langle \vartheta, \pi H\mu_\beta^V \rangle_{L^2}.$$

This shows that $U^{\mu_\beta^V}$ has a distributional derivative given by $\pi H\mu_\beta^V$. Moreover, since the density $\mu_\beta^V(x)$ belongs to $L^2(\mathbb{T})$ by Proposition 2.1 (b), so does $H\mu_\beta^V$ and thus $(U^{\mu_\beta^V})' \in L^2(\mathbb{T})$. \square

Proposition 2.5. *If $V \in H^1(\mathbb{T})$ then $U^{\mu_\beta^V} \in \mathcal{C}^{1,1/2}(\mathbb{T})$. Moreover, if $V \in \mathcal{C}^{m,1}(\mathbb{T})$ for some $m \geq 0$, then $\mu_\beta^V \in \mathcal{C}^{m,1}(\mathbb{T})$.*

Note that $V \in H^1(\mathbb{T})$ implies that V is continuous and this ensures the existence of μ_β^V .

Proof. By differentiating the Euler-Lagrange equation (16) we obtain the distributional identity

$$(\mu_\beta^V)' = -\mu_\beta^V(2\beta\pi H\mu_\beta^V + V'). \tag{21}$$

Since $H\mu_\beta^V \in L^2(\mathbb{T})$ and $\|\mu_\beta^V\|_{L^\infty} < \infty$ according to Proposition 2.1 (b), (21) yields that $\mu_\beta^V \in H^1(\mathbb{T})$ as soon as $V' \in L^2(\mathbb{T})$. This also shows that $(H\mu_\beta^V)' = H(\mu_\beta^V)' \in L^2(\mathbb{T})$ and thus $H\mu_\beta^V \in H^1(\mathbb{T}) \subset C^{1/2}(\mathbb{T})$. In particular, the first claim follows by Lemma 2.4. Moreover, if we further assume that $V \in C^{0,1}(\mathbb{T})$, then $\|(\mu_\beta^V)'\|_{L^\infty} < \infty$ by (21) and the second statement is proven for $m = 0$.

Next, we differentiate (21) in order to obtain

$$(\mu_\beta^V)'' = -(\mu_\beta^V)'(2\beta\pi H\mu_\beta^V + V') - \mu_\beta^V(2\beta\pi H(\mu_\beta^V)' + V''). \tag{22}$$

If we assume $V \in C^{1,1}(\mathbb{T})$, then in particular it is $C^{0,1}$ and we have already shown that $\|(\mu_\beta^V)'\|_{L^\infty} < \infty$. Together with (22) this provides $(\mu_\beta^V)'' \in L^2(\mathbb{T})$, thus $(H(\mu_\beta^V)')' = H(\mu_\beta^V)'' \in L^2(\mathbb{T})$, and this yields in turn $H(\mu_\beta^V)' \in C^{1/2}(\mathbb{T})$. Using (22) again, we obtain $\|(\mu_\beta^V)''\|_{L^\infty} < \infty$ and the claim holds for $m = 1$.

The case $m \geq 2$ follows inductively by differentiating (22) and using the same reasoning. \square

3. Proof of Theorem 1.4

We now turn to the proof of Theorem 1.4. The proof follows the same strategy as the one in [8] and is based on combining a Coulomb transport inequality together with an energy estimate after an appropriate regularization of the empirical measure. The regularization we use here is rather similar to [25] and the technical input with this respect here is the following lemma.

Lemma 3.1. *Given any configuration of distinct points $x_1, \dots, x_N \in \mathbb{T}$, there exists a configuration $y_1, \dots, y_N \in \mathbb{T}$ satisfying:*

$$\min_{j \neq k} |e^{iy_j} - e^{iy_k}| \geq \frac{2}{5N^4}, \quad \sum_{j=1}^N |e^{ix_j} - e^{iy_j}| \leq 1,$$

and

$$\sum_{j \neq k} \log \frac{1}{|e^{ix_j} - e^{ix_k}|} \geq \sum_{j \neq k} \log \frac{1}{|e^{iy_j} - e^{iy_k}|} - N.$$

Proof. Given any ordered configuration $x_1 < \dots < x_N$ in \mathbb{T} , there exists at least one index j such that $x_{j+1} - x_j \geq 2\pi/N$. Thus, by permutation and translation, one can assume without loss of generality that

$$-\pi + \frac{2}{N} \leq x_1 < \dots < x_N < \pi - \frac{2}{N}.$$

Consider the increasing bijection $x \in \mathbb{T} \mapsto \tilde{x} := \tan(x/2) \in \mathbb{R} \cup \{\pm\infty\}$ which satisfies

$$|e^{ix} - e^{iy}| = \frac{2|\tilde{x} - \tilde{y}|}{\sqrt{1 + \tilde{x}^2}\sqrt{1 + \tilde{y}^2}}.$$

We set $\tilde{y}_1 := \tilde{x}_1$ and $\tilde{y}_{j+1} := \tilde{y}_j + \max(\tilde{x}_{j+1} - \tilde{x}_j, N^{-2})$ and then let $y_1 < \dots < y_N \in \mathbb{T}$ be the configuration obtained by taking the image of the \tilde{y}_j 's by the inverse bijection. Since by construction $\tilde{y}_j - \tilde{x}_j \leq (j - 1)N^{-2}$ we have

$$\sum_{j=1}^N |e^{ix_j} - e^{iy_j}| \leq 2N^{-2} \sum_{j=1}^N (j - 1) \leq 1.$$

Next, by assumption on the x_j 's we have $\max_j |\tilde{x}_j| \leq |\tan(\frac{\pi}{2} - \frac{1}{N})| \leq N$, which yields $\max_j |\tilde{y}_j| \leq 2N$, and we thus obtain, for every $j \neq k$,

$$|e^{iy_j} - e^{iy_k}| \geq \frac{2}{N^2(1 + 4N^2)} \geq \frac{2}{5N^4}.$$

Finally, we have

$$\begin{aligned} \sum_{j \neq k} \log \frac{2}{|e^{ix_j} - e^{ix_k}|} &= \sum_{j \neq k} \log \frac{1}{|\tilde{x}_j - \tilde{x}_k|} + (N - 1) \sum_{j=1}^N \log(1 + \tilde{x}_j^2) \\ &\geq \sum_{j \neq k} \log \frac{1}{|\tilde{y}_j - \tilde{y}_k|} + (N - 1) \sum_{j=1}^N \log(1 + \tilde{x}_j^2) \\ &= \sum_{j \neq k} \log \frac{2}{|e^{iy_j} - e^{iy_k}|} + (N - 1) \sum_{j=1}^N \log \left(\frac{1 + \tilde{x}_j^2}{1 + \tilde{y}_j^2} \right) \\ &\geq \sum_{j \neq k} \log \frac{2}{|e^{iy_j} - e^{iy_k}|} - (N - 1) \sum_{j=1}^N \log \left(\frac{1 + (\tilde{x}_j + N^{-2}(j - 1))^2}{1 + \tilde{x}_j^2} \right). \end{aligned}$$

Using that, for any $0 < c < 1$,

$$\max_{x \in \mathbb{R}} \log \left(\frac{1 + (x + c)^2}{1 + x^2} \right) = \log \left(1 + \frac{2c}{\sqrt{c^2 + 4} - c} \right) \leq 2c,$$

we obtain

$$(N - 1) \sum_{j=1}^N \log \left(\frac{1 + (x_j + N^{-2}(j - 1))^2}{1 + x_j^2} \right) \leq N$$

which completes the proof of the lemma. \square

Proof of Theorem 1.4. Recalling (4), if we set for convenience

$$g(x) := \log \left| \sin \left(\frac{x}{2} \right) \right|^{-1}$$

then we can write

$$d\mathbb{P}_N(x_1, \dots, x_N) = \frac{1}{Z'_N} \exp \left\{ -\frac{\beta}{N} \sum_{j \neq k} g(x_j - x_k) - \sum_{j=1}^N V(x_j) \right\} \prod_{j=1}^N \frac{dx_j}{2\pi} \tag{23}$$

for some new normalization constant $Z'_N > 0$.

Step 1: Lower bound on the partition function. By writing

$$Z'_N = \int \exp \left\{ -\frac{\beta}{N} \sum_{j \neq k} g(x_j - x_k) - \sum_{j=1}^N (V(x_j) + \log(2\pi\mu_\beta^V(x_j))) \right\} \prod_{j=1}^N \mu_\beta^V(dx_j)$$

and using Jensen’s inequality, we obtain

$$\begin{aligned} \log Z'_N &\geq \int \left(-\frac{\beta}{N} \sum_{j \neq k} g(x_j - x_k) - \sum_{j=1}^N (V(x_j) + \log(2\pi\mu_\beta^V(x_j))) \right) \prod_{j=1}^N \mu_\beta^V(dx_j) \\ &= -\beta(N - 1)\mathcal{E}(\mu_\beta^V) - N \int (V + \log(2\pi\mu_\beta^V)) d\mu_\beta^V \\ &= -NF_\beta(\mu_\beta^V) + \beta\mathcal{E}(\mu_\beta^V). \end{aligned} \tag{24}$$

Step 2: Regularization and energy estimates. Given any configuration $x_1, \dots, x_N \in \mathbb{T}$ of distinct points, let $y_1, \dots, y_N \in \mathbb{T}$ be as in Lemma 3.1 and set

$$\tilde{\mu}_N := \frac{1}{N} \sum_{i=1}^n \delta_{y_i} * \lambda_{N^{-5}}, \quad \lambda_\varepsilon := \mathbf{1}_{[0,\varepsilon]}(x) \frac{dx}{\varepsilon}.$$

Since $g'(x) = -(2 \tan(x/2))^{-1}$, a Taylor-Lagrange expansion yields for any $|u| \leq |x|/2$,

$$|g(x + u) - g(x)| \leq \frac{|u|}{2 \sin |x/4|} \leq \frac{|u|}{\sin |x/2|}. \tag{25}$$

Since $\sin(|y_j - y_k|/2) \geq N^{-4}/5$ and $\sin(|y_k - y_j - u|/2) \geq N^{-4}/5 - |u|/2$ when $j \neq k$, we deduce from Lemma 3.1 and (25) that for any $N \geq 10$,

$$\begin{aligned}
 \sum_{j \neq k} g(x_j - x_k) &\geq \sum_{j \neq k} g(y_j - y_k) - N \\
 &\geq \sum_{j \neq k} \int g(y_j - y_k + u) \lambda_{N-5}(du) - 7N/2 \\
 &\geq \sum_{j \neq k} \iint g(y_j - y_k + u - v) \lambda_{N-5}(du) \lambda_{N-5}(dv) - 7N \\
 &= N^2 \mathcal{E}(\tilde{\mu}_N) - N \mathcal{E}(\lambda_{N-5}) - 7N.
 \end{aligned} \tag{26}$$

Next, by using that $2|\sin(a)| \geq |a|$ when $|a| \leq 1$, we obtain the upper bound

$$\mathcal{E}(\lambda_\varepsilon) \leq \int \log \frac{1}{|x-y|} \lambda_\varepsilon(dx) \lambda_\varepsilon(dy) + \log 2 = -\log \varepsilon + 3/2 + \log 2. \tag{27}$$

If we set $c := \mathcal{E}(\mu_\beta^V) + 17/2 + \log 2$, then by combining (24)–(27) we obtain,

$$\begin{aligned}
 &\frac{1}{Z'_N} \exp \left\{ -\frac{\beta}{N} \sum_{j \neq k} g(x_j - x_k) - \sum_{j=1}^n V(x_j) \right\} \\
 &\leq e^{\beta(5 \log N + c)} e^{-N(\beta \mathcal{E}(\tilde{\mu}_N) - F_\beta(\mu_\beta^V)) - \sum_{j=1}^N V(x_j)} \\
 &= e^{\beta(5 \log N + c)} e^{-N(F_\beta(\tilde{\mu}_N) - F_\beta(\mu_\beta^V) - \mathcal{K}(\tilde{\mu}_N | \mu_\beta^V)) + N \int Q d(\tilde{\mu}_N - \mu_N)} \prod_{j=1}^N 2\pi \mu_\beta^V(x_j)
 \end{aligned}$$

where $Q := V + \log \mu_\beta^V = C_\beta^V - 2\beta U \mu_\beta^V$ by (16). Using Corollary 2.2, we deduce from (23) that for any $r > 0$,

$$\mathbb{P}_N(\mathcal{E}(\tilde{\mu}_N - \mu_\beta^V) > r) \leq e^{-\beta(Nr - 5 \log N - c)} \int e^{N \int Q d(\tilde{\mu}_N - \mu_N)} (\mu_\beta^V)^{\otimes N}(dx). \tag{28}$$

Finally, since by assumption $V \in H^1(\mathbb{T})$, Proposition 2.5 yields $U \mu_\beta^V$ is Lipschitz and, using again Lemma 3.1, we have

$$\begin{aligned}
 \left| N \int Q d(\tilde{\mu}_N - \mu_N) \right| &\leq 2\beta \sum_{j=1}^N \int |U \mu_\beta^V(y_j + u) - U \mu_\beta^V(x_j)| \lambda_{N-5}(du) \\
 &\leq 2\beta \|U \mu_\beta^V\|_{\text{Lip}} \left(\sum_{j=1}^N |e^{ix_j} - e^{iy_j}| + 2N \int |\sin(u/2)| \lambda_{N-5}(du) \right) \\
 &\leq 3\beta \|U \mu_\beta^V\|_{\text{Lip}}.
 \end{aligned}$$

Together with (28), we have finally obtained the energy estimate

$$\mathbb{P}_N(\mathcal{E}(\tilde{\mu}_N - \mu_\beta^V) > r) \leq e^{-\beta(Nr - 5 \log N - \tilde{C})}, \tag{29}$$

where $\tilde{C} := \mathcal{E}(\mu_\beta^V) + 3\|U^{\mu_\beta^V}\|_{\text{Lip}} + 17/2 + \log 2$.

Step 4: The Coulomb transport inequality and conclusion. Lemma 3.1 yields,

$$\begin{aligned} W_1(\mu_N, \tilde{\mu}_N) &\leq \frac{1}{N} \sum_{j=1}^N \int |e^{ix_j} - e^{i(y_j+u)}| \lambda_{N-5}(du) \\ &\leq \frac{1}{N} + 2 \int |\sin(u/2)| \lambda_{N-5}(du) \leq \frac{2}{N}. \end{aligned}$$

Since both $\tilde{\mu}_N$ and μ_β^V have finite logarithmic energy, it follows from [8, Theorem 1.1] and the discussion below that, for every $\varepsilon > 0$,

$$W_1(\tilde{\mu}_N, \mu_\beta^V)^2 \leq 4\pi \mathcal{E}(\tilde{\mu}_N - \mu_\beta^V).$$

Moreover, using that

$$\frac{1}{2} W_1(\mu_N, \mu_\beta^V)^2 \leq W_1(\tilde{\mu}_N, \mu_\beta^V)^2 + W_1(\mu_N, \tilde{\mu}_N)^2 \leq W_1(\tilde{\mu}_N, \mu_\beta^V)^2 + \frac{4}{N^2},$$

we obtain for any $r > 0$ from (29),

$$\begin{aligned} \mathbb{P}_N\left(W_1(\mu_N, \mu_\beta^V) > r\right) &\leq \mathbb{P}_N\left(W_1(\tilde{\mu}_N, \mu_\beta^V)^2 > \frac{r^2}{2} - \frac{4}{N^2}\right) \\ &\leq \mathbb{P}_N\left(\mathcal{E}(\tilde{\mu}_N - \mu_\beta^V) > \frac{1}{C_T}\left(\frac{r^2}{2} - \frac{4}{N^2}\right)\right) \\ &\leq e^{-\beta\left(\frac{1}{8\pi}Nr^2 - 5\log N - C\right)} \end{aligned}$$

where the constant is given by

$$C := \mathcal{E}(\mu_\beta^V) + 3\|U^{\mu_\beta^V}\|_{\text{Lip}} + \frac{17}{2} + \log 2 + \frac{1}{\pi} \tag{30}$$

and the proof of the theorem is complete. \square

4. Main steps for the proof of Theorem 1.2

In this section, we explain the main strategy to prove Theorem 1.2. It is based on the multi-dimensional Gaussian approximation result from [24] combined with the previous concentration inequality and a study of the spectral properties of the operator \mathcal{L} .

Consider the differential operator given by

$$\begin{aligned} \mathbf{L} &:= \Delta - \nabla\left(\frac{2\beta}{N}\mathcal{H}(\mathbf{x}) + \sum_{j=1}^N V(x_j)\right) \cdot \nabla \\ &= \sum_{j=1}^N \partial_{x_j}^2 + \frac{2\beta}{N} \sum_{i \neq j} \frac{\partial_{x_j}}{2 \tan\left(\frac{x_j - x_i}{2}\right)} - \sum_{j=1}^N V'(x_j) \partial_{x_j}, \end{aligned}$$

which satisfies the integration by part identity $\int f(-\mathbf{L}g) d\mathbb{P}_N = \int \nabla f \cdot \nabla g d\mathbb{P}_N$ for any smooth functions $f, g : \mathbb{T}^N \rightarrow \mathbb{R}$.

Recalling that $\nu_N = \sqrt{N}(\hat{\mu}_N - \mu_\beta^V)$, we first show that $\nu_N(\phi)$, seen as a mapping $\mathbb{T}^N \rightarrow \mathbb{R}$, is an *approximate eigenfunction* for \mathbf{L} as long as ϕ is a (strong) eigenfunction of the differential operator \mathcal{L} defined in (10). More precisely, we have the approximate commutation relation:

Lemma 4.1. *For any $\phi \in C^2(\mathbb{T})$ we have*

$$\mathbf{L} \nu_N(\phi) = -\nu_N(\mathcal{L}\phi) + \frac{\beta}{\sqrt{N}}\zeta_N(\phi)$$

where we introduced

$$\zeta_N(\phi) := \iint \frac{\phi'(x) - \phi'(y)}{2 \tan(\frac{x-y}{2})} \nu_N(dx)\nu_N(dy) - \int \phi'' d\mu_N. \tag{31}$$

Proof. If we set $\Phi(\mathbf{x}) := \sum_{j=1}^N \phi(x_j)$ then we have

$$\begin{aligned} \mathbf{L}\Phi(\mathbf{x}) &= \sum_{j=1}^N \phi''(x_j) + \frac{\beta}{N} \sum_{i \neq j}^N \frac{\phi'(x_j)}{\tan(\frac{x_j-x_i}{2})} - \sum_{j=1}^N \phi'(x_j)V'(x_j) \\ &= \left(1 - \frac{\beta}{N}\right) \sum_{j=1}^N \phi''(x_j) + \frac{\beta}{N} \sum_{i,j=1}^N \frac{\phi'(x_j) - \phi'(x_i)}{2 \tan(\frac{x_j-x_i}{2})} - \sum_{j=1}^N \phi'(x_j)V'(x_j). \end{aligned} \tag{32}$$

Next, it is convenient to introduce the operator Ξ defined by

$$\Xi\psi(x) := \int \frac{\psi(x) - \psi(t)}{2 \tan(\frac{x-t}{2})} \mu_\beta^V(dt), \tag{33}$$

which is a weighted version of the Hilbert transform H defined in (11). Indeed, we can write

$$\begin{aligned} &\frac{1}{N} \sum_{i,j=1}^N \frac{\phi'(x_j) - \phi'(x_i)}{2 \tan(\frac{x_j-x_i}{2})} \\ &= 2\sqrt{N} \int \Xi(\phi') d\nu_N + N \int \Xi(\phi') d\mu_\beta^V + \iint \frac{\phi'(x) - \phi'(y)}{2 \tan(\frac{x-y}{2})} \nu_N(dx)\nu_N(dy) \end{aligned}$$

and this yields together with (32) and (31),

$$\mathbf{L}\Phi = N \int (\phi'' + \beta\Xi(\phi') - \phi'V') d\mu_\beta^V + \sqrt{N} \int (\phi'' + 2\beta\Xi(\phi') - \phi'V') d\nu_N + \beta\zeta_N(\phi). \tag{34}$$

By (10), observe that the variational equation (21) yields

$$\phi'' + 2\beta \Xi(\phi') - V'\phi' = -\mathcal{L}\phi \tag{35}$$

where we used that by (33),

$$\Xi\psi = \pi(H(\psi\mu_\beta^V) - \psi H(\mu_\beta^V)).$$

Moreover, we obtain by using that $H^* = -H$,

$$\begin{aligned} \int \Xi\psi \, d\mu_\beta^V &= -\pi \int H(\mu_\beta^V)\psi \, d\mu_\beta^V + \pi \langle H(\mu_\beta^V\psi), \mu_\beta^V \rangle_{L^2} \\ &= -2\pi \int H(\mu_\beta^V)\psi \, d\mu_\beta^V. \end{aligned}$$

By integrating (21) against $\phi' \, dx$, this yields together with an integration by parts:

$$\int (\phi'' + \beta \Xi(\phi') - \phi'V') \, d\mu_\beta^V = 0. \tag{36}$$

By combining (34)–(36), we have finally shown that

$$\mathbf{L}\Phi = -\sqrt{N}\nu_N(\mathcal{L}\phi) + \beta\zeta_N(\phi)$$

and the result follows by linearity of \mathbf{L} since $\Phi = \sqrt{N}\nu_N(\phi) + N \int \phi \, d\mu_\beta^V$. \square

It turns out the random variables $\zeta_N(\phi)$ are of smaller order of magnitude than the fluctuations provided ϕ is smooth enough. More precisely, we have the following estimates.

Lemma 4.2. *There exists a constant $C = C(\beta, V) > 0$ such that, for any $N \geq 10$, for any $\phi \in \mathcal{C}^{3,1}(\mathbb{T})$, we have*

$$\mathbb{E} \left[|\zeta_N(\phi)|^2 \right] \leq C \|\phi'''\|_{\text{Lip}}^2 (\log N)^2.$$

Moreover, for any Lipschitz function $g : \mathbb{T} \rightarrow \mathbb{R}$

$$\mathbb{E} \left[\left| \int g \, d\nu_N \right|^2 \right] \leq C \|g\|_{\text{Lip}}^2 \log N.$$

The proof of this lemma is based on Theorem 1.4 and is postponed to Section 5.

Another important input is the existence of an eigenbasis of \mathbf{H} for the operator \mathcal{L} that behaves like an eigenbasis of a Sturm-Liouville operator. Note that by (12), \mathbf{H} is a

separable Hilbert space and it follows from Proposition 2.1(b) that the associated norm satisfies

$$\delta \|\psi'\|_{L^2}^2 \leq \|\psi\|_{\mathbf{H}}^2 \leq \delta^{-1} \|\psi'\|_{L^2}^2,$$

a fact we will use at several instances below.

Proposition 4.3. *Assume that $V \in C^{m,1}(\mathbb{T})$ for some $m \geq 1$. Then there exists a family $(\phi_j)_{j=1}^\infty$ of functions $\phi_j : \mathbb{T} \rightarrow \mathbb{R}$ such that:*

- (a) $\mathcal{L}\phi_j = \varkappa_j \phi_j$ where $(\varkappa_j)_{j=1}^\infty$ is an increasing sequence of positive numbers.
- (b) $(\phi_j)_{j=1}^\infty$ is an orthonormal basis of the Hilbert space \mathbf{H} .
- (c) There exists $\alpha > 0$ such that $\varkappa_j \sim \alpha j^2$ as $j \rightarrow \infty$.
- (d) $\phi_j \in C^m(\mathbb{T})$ and there exist constants C_k such that for every $j \geq 1$ and $k \in \{0, \dots, m\}$,

$$\|\phi_j^{(k)}\|_{\text{Lip}} \leq C_k \varkappa_j^{\frac{k+1}{2}}.$$

The proof of Proposition 4.3 is postponed to the sections 6 and 7.

Proposition 4.4. *Assume that the external potential $V \in C^{3,1}(\mathbb{T})$. There exists a constant $C = C(\beta, V) > 0$ such that, if we set*

$$F := (\sqrt{\varkappa_1} \nu_N(\phi_1), \dots, \sqrt{\varkappa_d} \nu_N(\phi_d)),$$

then we have for every $N \geq 10$ and $d \geq 1$,

$$W_2(F, \mathcal{N}(0, I_d)) \leq \frac{C}{\sqrt{N}} \left(\beta \log N \sqrt{\sum_{j=1}^d \varkappa_j^3} + \sqrt{\log N \sum_{j=1}^d \varkappa_j^2 \sum_{j=1}^d \varkappa_j} \right).$$

Here $\mathcal{N}(0, I_d)$ stands for a real standard Gaussian random vector in \mathbb{R}^d . This proposition is a consequence of the previous concentration estimates together with the following general normal approximation given by [24, Proposition 2.1]; for $F = (F_1, \dots, F_d) \in C^2(\mathbb{T}^N, \mathbb{R}^d)$ we set $\mathbf{L}F := (\mathbf{L}F_1, \dots, \mathbf{L}F_d)$ and denote by $\|\cdot\|_{\mathbb{R}^d}$ the Euclidean norm of \mathbb{R}^d .

Theorem 4.5. *For any given $F = (F_1, \dots, F_d) \in C^2(\mathbb{T}^N, \mathbb{R}^d)$, let*

$$\Gamma := \left[\nabla F_i \cdot \nabla F_j \right]_{i,j=1}^d$$

and see both F and Γ as random variables defined on the probability space $(\mathbb{T}^N, \mathfrak{B}(\mathbb{T})^{\otimes N}, \mathbb{P}_N)$. Given any $d \times d$ diagonal matrix K with positive diagonal entries, we have

$$W_2(F, \mathcal{N}(0, \mathbf{I}_d)) \leq \sqrt{\mathbb{E} \left[\|F + K^{-1} \mathbf{L}F\|_{\mathbb{R}^d}^2 \right]} + \sqrt{\mathbb{E} \left[\|\mathbf{I}_d - K^{-1} \Gamma\|_{\mathbb{R}^{d \times d}}^2 \right]}.$$

Proof of Proposition 4.4. By Proposition 4.3 (a) and Lemma 4.1, we have for every $j \geq 1$,

$$\mathbf{L}F_j = -\varkappa_j F_j + \beta \sqrt{\frac{\varkappa_j}{N}} \zeta_N(\phi_j).$$

As a consequence, taking $K := \text{diag}(\varkappa_1, \dots, \varkappa_d)$, we obtain

$$\mathbb{E} \left[\|F + K^{-1} \mathbf{L}F\|_{\mathbb{R}^d}^2 \right] = \frac{\beta^2}{N} \sum_{j=1}^d \varkappa_j^{-1} \mathbb{E} \left[|\zeta_N(\phi_j)|^2 \right]$$

and, by Lemma 4.2 and Proposition 4.3 (c)-(d), this yields

$$\mathbb{E} \left[\|F + K^{-1} \mathbf{L}F\|_{\mathbb{R}^d}^2 \right] \leq C\beta^2 \frac{(\log N)^2}{N} \sum_{j=1}^d \varkappa_j^3. \tag{37}$$

Next, since for any $i, j \in \mathbb{N}$,

$$\Gamma_{ij} = \nabla F_i \cdot \nabla F_j = \sqrt{\varkappa_i \varkappa_j} \int \phi'_i \phi'_j \, d\mu_N,$$

and using that the ϕ_j 's are orthonormal, we obtain

$$\begin{aligned} \mathbb{E} \left[\|\mathbf{I}_d - K^{-1} \Gamma\|_{\mathbb{R}^{d \times d}}^2 \right] &= \mathbb{E} \left[\|\mathbf{I}_d - K^{-1/2} \Gamma K^{-1/2}\|_{\mathbb{R}^{d \times d}}^2 \right] \\ &= \frac{1}{N} \sum_{i,j=1}^d \mathbb{E} \left[\left| \int \phi'_i \phi'_j \, d\nu_N \right|^2 \right]. \end{aligned}$$

Proposition 4.3 (d) moreover yields, for any $i, j \in \mathbb{N}$,

$$\|\phi'_i \phi'_j\|_{\text{Lip}} \leq \|\phi'_i\|_{\text{Lip}} \|\phi_j\|_{\text{Lip}} + \|\phi_i\|_{\text{Lip}} \|\phi'_j\|_{\text{Lip}} \leq C(\varkappa_i \sqrt{\varkappa_j} + \sqrt{\varkappa_i} \varkappa_j)$$

and it thus follows from Lemma 4.2 that

$$\mathbb{E} \left[\|\mathbf{I}_d - K^{-1} \Gamma\|_{\mathbb{R}^{d \times d}}^2 \right] \leq C \frac{\log N}{N} \sum_{j=1}^d \varkappa_j^2 \sum_{j=1}^d \varkappa_j. \tag{38}$$

The proposition follows by combining estimates (37) and (38) together with Theorem 4.5. \square

We are finally in position to prove Theorem 1.2 by decomposing a general test function into the eigenbasis $(\phi_j)_{j=1}^\infty$ and by using Proposition 4.4.

Proof of Theorem 1.2. Assume that $V \in \mathcal{C}^{3,1}(\mathbb{T})$ and let $\psi \in \mathcal{C}^{2\gamma+1}(\mathbb{T})$ for some integer $\gamma \geq 2$. We can assume without loss of generality that $\int \psi \, d\mu_\beta^V = 0$. Thus $\psi \in \mathbb{H}$ and we have by Proposition 4.3 (b),

$$\psi \stackrel{\mathbb{H}}{=} \sum_{j=1}^{\infty} \langle \psi, \phi_j \rangle_{\mathbb{H}} \phi_j \tag{39}$$

Moreover, since ψ lies in the domain of \mathcal{L}^γ and using that \mathcal{L} is symmetric, we have

$$|\langle \psi, \phi_j \rangle_{\mathbb{H}}| = \frac{1}{\varkappa_j^\gamma} |\langle \mathcal{L}^\gamma \psi, \phi_j \rangle_{\mathbb{H}}| \leq \frac{1}{\varkappa_j^\gamma} \|\mathcal{L}^\gamma \psi\|_{\mathbb{H}}. \tag{40}$$

In particular, by Proposition 4.3 (c), the series (39) converges uniformly on \mathbb{T} .

Next, given any $d \in \mathbb{N}$, let us consider the truncation of ψ ,

$$\psi^{[d]} := \sum_{j=1}^d \langle \psi, \phi_j \rangle_{\mathbb{H}} \phi_j.$$

Proposition 4.3 (c)-(d) and the upper bound (40) yield for $\gamma > 1$,

$$\begin{aligned} \|\psi - \psi^{[d]}\|_{\text{Lip}} &\leq C \sum_{j=d+1}^{\infty} |\langle \psi, \phi_j \rangle_{\mathbb{H}}| \sqrt{\varkappa_j} \\ &\leq C \|\mathcal{L}^\gamma \psi\|_{\mathbb{H}} \sum_{j=d+1}^{\infty} \varkappa_j^{-\gamma+1/2} \\ &\leq C \|\mathcal{L}^\gamma \psi\|_{\mathbb{H}} d^{-2(\gamma-1)}. \end{aligned}$$

Thus, by definition of the W_2 metric and Lemma 4.2, this yields

$$W_2 \left(\nu_N(\psi), \nu_N(\psi^{[d]}) \right) \leq \sqrt{\mathbb{E} \left[\left| \int (\psi - \psi^{[d]}) \, d\nu_N \right|^2 \right]} \leq C \|\mathcal{L}^\gamma \psi\|_{\mathbb{H}} \frac{\sqrt{\log N}}{d^{2(\gamma-1)}}. \tag{41}$$

Next, if we set $\eta_m := \sum_{j=1}^m \varkappa_j^{-1} |\langle \psi, \phi_j \rangle_{\mathbb{H}}|^2$ for $m \in \mathbb{N} \cup \{\infty\}$, then we obtain from (40) and Proposition 4.3 (c) that

$$\begin{aligned} W_2(\mathcal{N}(0, \eta_d), \mathcal{N}(0, \eta_\infty))^2 &= (\sqrt{\eta_\infty} - \sqrt{\eta_d})^2 \\ &\leq \eta_\infty - \eta_d \\ &= \sum_{j=d+1}^{\infty} \varkappa_j^{-1} |\langle \psi, \phi_j \rangle_{\mathbb{H}}|^2 \\ &\leq C \|\mathcal{L}^\gamma \psi\|_{\mathbb{H}}^2 d^{-4\gamma-1}. \end{aligned} \tag{42}$$

Last, Proposition 4.4 and Proposition 4.3 (c) yield

$$\begin{aligned}
 W_2\left(\nu_N(\psi^{[d]}), \mathcal{N}(0, \eta_d)\right) &= W_2\left(\sum_{j=1}^d \frac{\langle \psi, \phi_j \rangle_{\mathbb{H}}}{\sqrt{\kappa_j}} \nu_N(\phi_j), \mathcal{N}(0, \eta_d)\right) \\
 &\leq \sqrt{\eta_\infty} W_2(F, \mathcal{N}(0, \mathbf{I}_d)) \\
 &\leq \sqrt{\frac{C\eta_\infty}{N}} \left(\beta \log N d^{7/2} + \sqrt{\log N} d^4\right). \tag{43}
 \end{aligned}$$

Finally, by combining the estimates (41)–(43) and taking for d the integer part of $N^{\frac{1/4}{\gamma+1}}$, we obtain

$$W_2(\nu_N(\psi), \mathcal{N}(0, \eta_\infty)) \leq C_\psi \sqrt{\frac{\log N}{N^{\frac{\gamma-1}{\gamma+1}}}}$$

where $C_\psi > 0$ depends on $\eta_\infty, \|\mathcal{L}^\gamma \psi\|_{\mathbb{H}}, \beta$ and V only. It remains to check that η_∞ equals to the variance $\sigma_\beta^V(\psi)^2$ given in (13); this is proven in Proposition 6.3 below. The proof of the theorem is therefore complete. \square

5. Concentration estimates: proof of Lemma 4.2

If we use the Kantorovich-Rubinstein dual representation of W_1 and take $r := R\sqrt{\log N/N}$ in Theorem 1.4, then under the same assumptions and using the same notation as in that theorem we obtain the following estimate: there exist $C = C(\mu_\beta^V, \beta) > 0$ and $\kappa = \kappa(\beta) > 0$ such that, for every $R \geq 6$ and $N \geq 10$,

$$\mathbb{P}_N \left(\sup_{\|f\|_{\text{Lip}} \leq 1} \left| \int f \, d\nu_N \right| > \sqrt{R \log N} \right) \leq CN^{-\kappa R}. \tag{44}$$

We also need the next estimate.

Lemma 5.1. *There exist $\kappa' = \kappa'(\beta) > 0$ and a constant $R_0 > 0$ such that, for any function $\psi \in \mathcal{C}^{2,1}(\mathbb{T})$, one has for every $R \geq R_0$ and $N \geq 10$,*

$$\mathbb{P}_N \left(\left| \iint \frac{\psi(x) - \psi(y)}{2 \tan(\frac{x-y}{2})} \nu_N(dx) \nu_N(dy) \right| > R \|\psi''\|_{\text{Lip}} \log N \right) \leq CN^{-\kappa' R}.$$

Proof. The strategy is to prove that the random function

$$\Psi_N(x) := \int \frac{\psi(x) - \psi(y)}{2 \tan(\frac{x-y}{2})} \nu_N(dy)$$

has Lipschitz constant controlled by $\|\psi''\|_{\text{Lip}}\sqrt{\log N}$ with high probability and then to use (44).

Since $\psi \in \mathcal{C}^{2,1}(\mathbb{T})$, we verify that for any $x \in \mathbb{R}$,

$$\Psi'_N(x) = - \int \frac{\psi(x) - \psi(y) - \psi'(x) \sin(x - y)}{4 \sin^2(\frac{x-y}{2})} \nu_N(dy).$$

We now provide an upper bound on the Lipschitz constant of the integrand of Ψ'_N which is uniform in x . Indeed, we have

$$\begin{aligned} & \frac{d}{dy} \left(\frac{\psi(x) - \psi(y) - \psi'(x) \sin(x - y)}{4 \sin^2(\frac{x-y}{2})} \right) \\ &= \frac{(\psi(x) - \psi(y)) \cos(\frac{x-y}{2}) - (\psi'(x) + \psi'(y)) \sin(\frac{x-y}{2})}{4 \sin^3(\frac{x-y}{2})}. \end{aligned} \tag{45}$$

Let us recall that we introduced the chordal metric $d_{\mathbb{T}}(x, y)$. Two Taylor-Lagrange expansions yield, for any $x, y \in \mathbb{T}$,

$$\psi(x) - \psi(y) = \frac{d_{\mathbb{T}}(x, y)}{2} (\psi'(x) + \psi'(y)) - \frac{d_{\mathbb{T}}(x, y)^2}{4} (\psi''(u) - \psi''(v))$$

for some $u, v \in \mathbb{T}$, so that

$$\psi(x) - \psi(y) = \frac{d_{\mathbb{T}}(x, y)}{2} (\psi'(x) + \psi'(y)) + \mathcal{O}(|x - y|^3) \|\psi''\|_{\text{Lip}}.$$

Together with (45), this implies that there exists a constant $C > 0$ such that

$$\sup_{x, y \in \mathbb{T}} \left| \frac{d}{dy} \left(\frac{\psi(x) - \psi(y) - \psi'(x) \sin(x - y)}{4 \sin^2(\frac{x-y}{2})} \right) \right| \leq C (\|\psi''\|_{\text{Lip}} + \|\psi'\|_{L^\infty}).$$

Since the mean value theorem yields that

$$\|\psi'\|_{L^\infty} \leq \pi \|\psi''\|_{L^\infty} \leq 2\pi \|\psi''\|_{\text{Lip}}, \tag{46}$$

we deduce from (44) that there exist constants $\kappa' > \kappa$ and $R_0 > 0$ such that for all $R \geq R_0$ and $N \geq 10$,

$$\mathbb{P}_N \left(\|\Psi_N\|_{\text{Lip}} > \|\psi''\|_{\text{Lip}} \sqrt{R \log N} \right) \leq CN^{-\kappa' R}.$$

Therefore, by (44) again, we obtain

$$\begin{aligned} & \mathbb{P}_N \left(\left| \iint \frac{\psi(x) - \psi(y)}{2 \tan(\frac{x-y}{2})} \nu_N(dx) \nu_N(dy) \right| > \|\psi''\|_{\text{Lip}} R \log N \right) \\ & \leq \mathbb{P}_N \left(\left| \int \Psi_N d\nu_N \right| > \|\psi''\|_{\text{Lip}} R \log N, \|\Psi_N\|_{\text{Lip}} \leq \|\psi''\|_{\text{Lip}} \sqrt{R \log N} \right) + CN^{-\kappa'R} \\ & \leq 2CN^{-\kappa'R}, \end{aligned}$$

which completes the proof of the lemma. \square

Proof of Lemma 4.2. Using (44) and that, for any real random variable X and $\alpha > 0$,

$$\mathbb{E}[X^2] = \alpha \int_0^\infty \mathbb{P}(|X| \geq \sqrt{\alpha R}) dR, \tag{47}$$

we obtain for any $N \geq 10$ and any Lipschitz function $g : \mathbb{T} \rightarrow \mathbb{R}$ that

$$\mathbb{E} \left[\left| \int g d\nu_N \right|^2 \right] \leq \left(6 + \frac{1}{\kappa \log N} \right) \|g\|_{\text{Lip}}^2 \log N$$

and the second statement of the lemma is obtained.

Next, by (31) and since μ_N is a probability measure, we have

$$|\zeta_N(\phi)| \leq \left| \iint \frac{\phi'(x) - \phi'(y)}{2 \tan(\frac{x-y}{2})} \nu_N(dx) \nu_N(dy) \right| + \|\phi''\|_{L^\infty}.$$

Using the inequality $\|\phi''\|_{L^\infty} \leq 2\pi\|\phi'''\|_{\text{Lip}}$ obtained as in (46), we deduce from Lemma 5.1 that for all $R \geq R_0$ and $N \geq 10$,

$$\mathbb{P} \left(|\zeta_N(\phi)| \geq 2R\|\phi'''\|_{\text{Lip}} \log N \right) \leq CN^{-\kappa'R}.$$

Thus, by (47), this yields

$$\mathbb{E} \left[|\zeta_N(\phi)|^2 \right] \leq 4\|\phi'''\|_{\text{Lip}}^2 (\log N)^2 \left(R_0^2 + \frac{2C}{(\kappa' \log N)^2} \right)$$

and the proof of the lemma is complete. \square

6. Spectral theory: proof of Proposition 4.3 (a)–(c)

In this section, we always assume $V \in \mathcal{C}^{1,1}(\mathbb{T})$. In particular it follows from Proposition 2.5 that $(\log \mu_\beta^V)'$ is Lipschitz continuous. Recalling (10), we write

$$\mathcal{L} = \mathcal{A} + 2\pi\beta \mathcal{W} \tag{48}$$

where we introduced the operators on $L^2(\mathbb{T})$,

$$\begin{aligned} \mathcal{A}\phi &:= -\phi'' - (\log \mu_\beta^V)' \phi' = -\frac{(\phi' \mu_\beta^V)'}{\mu_\beta^V}, \\ \mathcal{W}\phi &:= -H(\phi' \mu_\beta^V). \end{aligned} \tag{49}$$

Note that \mathcal{A} is a Sturm-Liouville operator in the sense that it reads

$$-\mathcal{A} = \frac{d}{dx} \left(p(x) \frac{d}{dx} \right) + q(x)$$

with $p := \log \mu_\beta^V$ and $q := 0$; we refer to [26,7] for general references on Sturm-Liouville equations.

We first check that \mathcal{L} is a positive operator on \mathbb{H} , as a consequence of the next lemma.

Lemma 6.1. *The operators \mathcal{A} and \mathcal{W} are both positive on \mathbb{H} .*

Proof. We have for any function $\phi \in \mathbb{H}$,

$$\langle \mathcal{A}\phi, \phi \rangle_{\mathbb{H}} = - \int_{\mathbb{T}} \left(\frac{\phi'}{\mu_\beta^V} \right)' \phi \, dx = \int_{\mathbb{T}} |\phi'|^2 \frac{dx}{\mu_\beta^V} \geq 0$$

where we set $\varphi := \phi' \mu_\beta^V$. Moreover, if one decomposes φ in the Fourier basis, then we have

$$\langle \mathcal{W}\phi, \phi \rangle_{\mathbb{H}} = - \int_{\mathbb{T}} H(\varphi)' \varphi \, dx = \sum_{k \in \mathbb{Z}} |k| |\hat{\varphi}_k|^2 = \|\varphi\|_{\mathbb{H}^{1/2}}^2 \geq 0 \tag{50}$$

and the lemma is proven. \square

The spectral properties of the Sturm-Liouville operator \mathcal{A} (with periodic boundary conditions) are well known, see for instance [7, Chapter 2 and 3], from which one can obtain the basic properties:

Lemma 6.2. *There exists an orthonormal basis $(\varphi_j)_{j=1}^\infty$ of \mathbb{H} consisting of (weak) eigenfunctions of \mathcal{A} associated with positive eigenvalues. Moreover, if*

$$\mathcal{A}\varphi_j = \lambda_j \varphi_j$$

with $0 < \lambda_1 \leq \lambda_2 \leq \dots$, then there exists $\alpha > 0$ such that, as $j \rightarrow \infty$,

$$\lambda_j \sim \alpha j^2.$$

Proof. Since for any smooth function $\phi : \mathbb{T} \rightarrow \mathbb{R}$ we have

$$\langle \mathcal{A}\phi, \phi \rangle_{L^2(\mu_\beta^V)} := \int \mathcal{A}\phi \phi \, d\mu_\beta^V = \int |\phi'|^2 \, d\mu_\beta^V \geq 0,$$

we see that \mathcal{A} is a positive Sturm-Liouville operator on $L^2(\mu_\beta^V)$ whose domain is $H^1(\mu_\beta^V) := \{\phi \in L^2(\mu_\beta^V) : \phi' \in L^2(\mu_\beta^V)\} = H^1(\mathbb{T})$, where we used Proposition 2.1 (b) for this equality. It then follows from the general properties of the Sturm-Liouville operators that there exists an orthonormal basis of $L^2(\mu_\beta^V)$ consisting of eigenfunctions $(\tilde{\varphi}_j)_{j=0}^{+\infty} \subset H^1(\mathbb{T})$ of \mathcal{A} associated to non-negative increasing eigenvalues $(\lambda_j)_{j=0}^\infty$. Moreover, by Weyl’s law (see e.g. [7, Theorem 3.3.2] in our setting), there exists $\alpha > 0$ such that $\lambda_j \sim \alpha j^2$ as $j \rightarrow \infty$.

The smallest eigenvalue $\lambda_0 = 0$ comes with the eigenfunction $\tilde{\varphi}_0 = 1$ which is orthogonal to H in $L^2(\mathbb{T})$, see (12). Since the $\tilde{\varphi}_j$ ’s are orthonormal in $L^2(\mu_\beta^V)$, we have for any $j \geq 1$,

$$\int \tilde{\varphi}_j \, d\mu_\beta^V = \langle \tilde{\varphi}_j, \tilde{\varphi}_0 \rangle_{L^2(\mu_\beta^V)} = 0 \tag{51}$$

and thus $(\tilde{\varphi}_j)_{j=1}^\infty \subset H$. Moreover, since we have for any $i, j \in \mathbb{N}$,

$$\begin{aligned} \langle \tilde{\varphi}_i, \tilde{\varphi}_j \rangle_H &= \langle \tilde{\varphi}'_i, \tilde{\varphi}'_j \rangle_{L^2(\mu_\beta^V)} \\ &= \langle \tilde{\varphi}_i, \mathcal{A}\tilde{\varphi}_j \rangle_{L^2(\mu_\beta^V)} \\ &= \lambda_j \delta_{ij}, \end{aligned}$$

it follows that $\lambda_1 > 0$ (since otherwise $\tilde{\varphi}_1$ would be a non-zero constant function and this would contradict (51)). Finally, if we set $\varphi_j := \tilde{\varphi}_j / \sqrt{\lambda_j}$, then the family $(\varphi_j)_{j=1}^\infty$ is an orthonormal basis of H that satisfies the requirements of the lemma. \square

Proposition 6.3. *Proposition 4.3 (a)–(c) hold true. More precisely, there exists an orthonormal basis $(\phi_j)_{j=1}^\infty$ of H such that $\mathcal{L}\phi_j = \varkappa_j \phi_j$ with $0 < \varkappa_1 \leq \varkappa_2 \leq \dots$ and we have*

$$\varkappa_j \sim \alpha j^2$$

as $j \rightarrow \infty$ for the same $\alpha > 0$ as in Lemma 6.2. In particular, \mathcal{L}^{-1} is a well defined trace class operator on H and, for any $\psi \in H$, we have

$$\langle \psi, \mathcal{L}^{-1}\psi \rangle_H = \sum_{j=1}^\infty \frac{1}{\varkappa_j} |\langle \psi, \phi_j \rangle_H|^2. \tag{52}$$

Proof. We use here basic results from operator theory, see e.g. [21]. Lemma 6.2 yields that \mathcal{A} is a positive self-adjoint operator on H and that \mathcal{A}^{-1} is trace-class. Since \mathcal{W} is

non-negative and self-adjoint on H , it follows that $\mathcal{L}^{-1} = (\mathcal{A} + 2\pi\beta\mathcal{W})^{-1}$ is a positive self-adjoint compact operator on H . The spectral theorem for self-adjoint compact operators then yields the existence of an orthonormal family $(\phi_j)_{j=1}^\infty$ in H and an increasing sequence of positive numbers $(\varkappa_j)_{j=1}^\infty$ such that $\mathcal{L}^{-1} = \sum_j \varkappa_j^{-1} \phi_j \otimes \phi_j$. In particular $\mathcal{L}\phi_j = \varkappa_j \phi_j$ weakly for every $j \geq 1$. Moreover, since \mathcal{L}^{-1} is positive, the family (ϕ_j) is necessarily a complete orthonormal family in H : part (a) and (b) are thus proven.

Writing $\langle \cdot, \cdot \rangle$ instead of $\langle \cdot, \cdot \rangle_H$ for simplicity, the min-max theorem (see e.g. [31, Theorem XIII.2]) yields, for any $j \geq 1$,

$$\varkappa_j = \max_{S_{j-1}} \min_{\substack{\psi \in S_{j-1}^\perp \\ \|\psi\|=1}} \langle \mathcal{L}\psi, \psi \rangle,$$

where the maximum is taken over all subspace $S_{j-1} \subset H$ of dimension $j - 1$. By taking $\tilde{S}_j := \text{span}(\varphi_1, \dots, \varphi_j)$ where $(\varphi_j)_{j=1}^\infty$ is as in Lemma 6.2, this provides

$$\begin{aligned} \varkappa_j &\geq \min_{\substack{\psi \in \tilde{S}_{j-1}^\perp \\ \|\psi\|=1}} \langle \mathcal{L}\psi, \psi \rangle \\ &\geq \min_{\substack{\psi \in \tilde{S}_{j-1}^\perp \\ \|\psi\|=1}} \langle \mathcal{A}\psi, \psi \rangle = \lambda_j \end{aligned} \tag{53}$$

where we also used that $\mathcal{W} \geq 0$ in the last inequality. Similarly, we use the reversed form of the min-max principle to obtain that, using also (48), for any $j \geq 1$,

$$\begin{aligned} \varkappa_j &= \min_{S_j} \max_{\substack{\psi \in S_j \\ \|\psi\|=1}} \langle \mathcal{L}\psi, \psi \rangle \\ &\leq \max_{\substack{\psi \in \tilde{S}_j \\ \|\psi\|=1}} \langle \mathcal{L}\psi, \psi \rangle \\ &\leq \lambda_j + 2\pi\beta \max_{\substack{\psi \in \tilde{S}_j \\ \|\psi\|=1}} \langle \mathcal{W}\psi, \psi \rangle. \end{aligned} \tag{54}$$

Next, by using (50), the Cauchy-Schwarz inequality, that H is an isometry of $L_0^2(\mathbb{T})$ such that $H(\psi)' = H(\psi')$ for every $\psi \in H^1(\mathbb{T})$, the second equality in (49) and Proposition 2.1 (b), we obtain for any $\ell \geq 1$,

$$\begin{aligned} \langle \mathcal{W}\varphi_\ell, \varphi_\ell \rangle_H &\leq 2\pi \|(\varphi'_\ell \mu_\beta^V)'\|_{L^2} \|\varphi'_\ell \mu_\beta^V\|_{L^2} \\ &= 2\pi \|\mu_\beta^V \mathcal{A}\varphi_\ell\|_{L^2} \|\varphi'_\ell \mu_\beta^V\|_{L^2} \\ &\leq 2\pi \delta^{-1} \|\mathcal{A}\varphi_\ell\|_{L^2(\mu_\beta^V)} \|\varphi_\ell\|_H \\ &= 2\pi \delta^{-1} \lambda_\ell^{1/2}. \end{aligned}$$

For the last step, we used that by definition, $\|\varphi_\ell\|_{\mathbf{H}} = 1$ and $\|\varphi_\ell\|_{L^2(\mu_\beta^V)} = \lambda_\ell^{-1/2}$ (see the end of the proof of Lemma 6.2). By (54), this yields

$$\varkappa_j \leq \lambda_j + 4\pi^2\beta\delta^{-1} \lambda_j^{1/2}.$$

Finally, combined with (53) and Lemma 6.2, the proof of the proposition is complete. \square

7. Regularity: proof of Proposition 4.3 (d)

We start with the following lemma.

Lemma 7.1. *Suppose that $V \in \mathcal{C}^{1,1}(\mathbb{T})$. There exists $C = C(\beta, V) > 0$ such that, if $\phi \in \mathbf{H}$ satisfies $\|\phi\|_{\mathbf{H}} = 1$ and $\mathcal{L}\phi = \varkappa\phi$ weakly for some $\varkappa > 0$, then $\phi'' \in L^2(\mathbb{T})$ and $\|\phi\|_{L^2} \leq C\varkappa^{-1/2}$.*

Proof. First, since $\int \phi d\mu_\beta^V = 0$ and ϕ is continuous, there exists $\xi \in \mathbb{T}$ such that $\phi(\xi) = 0$. Thus, by the Cauchy-Schwarz inequality,

$$\|\phi\|_{L^\infty} \leq \sup_{x \in \mathbb{T}} \left| \int \mathbf{1}_{[\xi, x]} \phi' d\theta \right| \leq 2\pi \|\phi'\|_{L^2}. \tag{55}$$

By Proposition 2.1(b), this yields in turn

$$\|\phi\|_{L^\infty} \leq 2\pi\delta^{-1/2} \|\phi\|_{\mathbf{H}} = 2\pi/\sqrt{\delta}. \tag{56}$$

Since $\mu_\beta^V \in \mathcal{C}^{1,1}(\mathbb{T})$ according to Proposition 2.5 and using that the Hilbert transform H preserves the $L^2(\mathbb{T})$ norm, we see the functions $H(\mu_\beta^V \phi')$ and $(\log \mu_\beta^V)' \phi'$ are in $L^2(\mathbb{T})$. Together with the definition (10) of \mathcal{L} , this implies that

$$-\phi'' = \varkappa\phi + 2\pi\beta H(\mu_\beta^V \phi') + (\log \mu_\beta^V)' \phi' \tag{57}$$

belongs to $L^2(\mathbb{T})$. Recalling (49), an integration by parts shows that

$$|\langle \phi, \mathcal{A}\phi \rangle_{L^2}| \leq \delta^{-1} \int \phi \mathcal{A}\phi d\mu_\beta^V = \delta^{-1} \|\phi\|_{\mathbf{H}}^2 = \delta^{-1}.$$

Moreover, by (49), using Cauchy-Schwarz inequality and (56), we have

$$\begin{aligned} |\langle \phi, \mathcal{W}\phi \rangle_{L^2}| &\leq \|\phi\|_{L^2} \|\phi' \mu_\beta^V\|_{L^2} \\ &\leq \delta^{-1/2} \|\phi\|_{L^\infty} \|\phi\|_{\mathbf{H}} \\ &\leq 2\pi\delta^{-1}. \end{aligned}$$

Put together, by (48), this yields

$$\|\phi\|_{L^2}^2 = \varkappa^{-1} \langle \phi, \mathcal{L}\phi \rangle_{L^2} \leq \varkappa^{-1} \delta^{-1} (1 + 4\pi^2\beta)$$

which completes the proof. \square

We finally turn to the proof of the last statement of Proposition 4.3 and thus complete the proof of Theorem 1.2.

Proof of Proposition 4.3 (d). Assume $V \in C^{m,1}(\mathbb{T})$ for some $m \geq 1$. In particular, Proposition 2.5 yields $\mu_\beta^V \in C^{m,1}(\mathbb{T})$ and, thanks to Proposition 2.1(b), we also have $\log \mu_\beta^V \in C^{m,1}(\mathbb{T})$ and thus $\|(\log \mu_\beta^V)^{(m+1)}\|_{L^\infty} < \infty$.

Starting from (57) and using that $\|H \cdot\|_{L^2} \leq \|\cdot\|_{L^2}$, that $\|\phi_j\|_{\mathbb{H}} = 1$ and Lemma 7.1, we see there exists $C = C(\beta, V) > 0$ such that, for any $j \geq 1$,

$$\begin{aligned} \|\phi_j''\|_{L^2} &\leq \varkappa_j \|\phi_j\|_{L^2} + 2\pi\beta \|\phi_j' \mu_\beta^V\|_{L^2} + \|(\log \mu_\beta^V)' \phi_j'\|_{L^2} \\ &\leq \varkappa_j \|\phi_j\|_{L^2} + (2\pi\beta\delta^{-1/2} + \delta^{-1/2} \|(\log \mu_\beta^V)'\|_{L^\infty}) \|\phi_j\|_{\mathbb{H}} \\ &\leq C\sqrt{\varkappa_j}. \end{aligned} \tag{58}$$

By (15) and (55) applied to ϕ_j' , this yields that Proposition 4.3 (d) holds true when $k = 0$.

Next, we use that for any $\psi \in H^1(\mathbb{T})$, we have by (55)

$$\|H\psi\|_{L^\infty} \leq 2\pi \|(H\psi)'\|_{L^2} = 2\pi \|\psi'\|_{L^2}.$$

Thus, since $\mu_\beta^V \in C^{1,1}(\mathbb{T})$ and $\phi_j'' \in L^2(\mathbb{T})$, according to (57), we have for every $j \geq 1$,

$$\begin{aligned} \|H(\mu_\beta^V \phi_j')\|_{L^\infty} &\leq 2\pi \|(\mu_\beta^V \phi_j')'\|_{L^2} \\ &\leq 2\pi (\|(\mu_\beta^V)'\|_{L^\infty} \|\phi_j'\|_{L^2} + \delta^{-1} \|\phi_j''\|_{L^2}) \\ &\leq C\sqrt{\varkappa_j} \end{aligned} \tag{59}$$

for some $C = C(\beta, V) > 0$; note that we used again that $\|\phi_j'\|_{L^2} \leq \delta^{-1/2} \|\phi_j\|_{\mathbb{H}}$. By using this estimate in (57) together with (56) and using the proposition for $k = 0$, we obtain

$$\|\phi_j''\|_{L^\infty} \leq C\varkappa_j.$$

This proves the proposition when $k = 1$. Note that, in particular, $\phi_j'' \in L^\infty(\mathbb{T})$.

Assume now that $m \geq 2$ so as to treat the case where $k = 2$. Observe that, since $\phi_j'' \in L^2(\mathbb{T})$, the right hand side of equation (57) has a weak derivative in L^2 and we obtain, for any $j \geq 1$,

$$-\phi_j''' = \varkappa_j \phi_j' + 2\pi\beta H(\mu_\beta^V \phi_j')' + (\log \mu_\beta^V)'' \phi_j' + (\log \mu_\beta^V)' \phi_j''. \tag{60}$$

Together with (58) and the upper bounds used to prove it, this yields

$$\|\phi_j'''\|_{L^2} \leq C\kappa_j$$

and in particular $\phi_j''' \in L^2(\mathbb{T})$. Similarly as in (59), this implies in turn that

$$\|H(\mu_\beta^V \phi_j')'\|_{L^\infty} \leq C\kappa_j.$$

By using this estimate combined together with the proposition for $k = 0$ and $k = 1$, we obtain from (60) that

$$\|\phi_j'''\|_{L^\infty} \leq C\kappa_j^{3/2}$$

and the proof of the proposition is complete when $k = 2$.

The setting where $k \geq 3$ is proven inductively by using the same method, after $k - 1$ differentiations of formula (57). \square

8. Continuity of the variance in the parameter $\beta \in [0, +\infty]$

In this final section, we study the limits of $\sigma_\beta^V(\psi)^2$ as $\beta \rightarrow 0$ and $\beta \rightarrow \infty$. We provide sufficient conditions on V so that the variance interpolates between the L^2 and the $H^{1/2}$ (semi-)norms, as it is the case when $V = 0$, see Lemma 1.3.

Convention: In this section, we denote the Hilbert space H and the operators \mathcal{L} , \mathcal{A} , and \mathcal{W} defined in the previous sections by H_β , \mathcal{L}_β , \mathcal{A}_β , and \mathcal{W}_β respectively to emphasize on the dependence on the parameter $\beta \geq 0$.

First, let us record the following smoothing property of the operators \mathcal{L}_β^{-1} .

Lemma 8.1. *Let $V \in C^{1,1}(\mathbb{T})$. If $f \in H_\beta$ for some $\beta > 0$, then $\mathcal{L}_\beta^{-1} f \in H^2(\mathbb{T})$.*

Proof. If $f \in H_\beta$, then by Proposition 4.3 we have the convergent expansion in H_β ,

$$\mathcal{L}_\beta^{-1} f = \sum_{j=1}^\infty \frac{\langle f, \phi_j \rangle_{H_\beta}}{\kappa_j} \phi_j.$$

By differentiating this formula and using the estimate (58) this shows that, if $V \in C^{1,1}(\mathbb{T})$, there exists a constant $C = C(\beta, V) > 0$ such that

$$\|(\mathcal{L}_\beta^{-1} f)''\|_{L^2} \leq \sum_{j=1}^\infty \frac{|\langle f, \phi_j \rangle_{H_\beta}|}{\kappa_j} \|\phi_j''\|_{L^2} \leq C \sum_{j=1}^\infty \frac{|\langle f, \phi_j \rangle_{H_\beta}|}{\sqrt{\kappa_j}} \leq C \|f\|_{H_\beta}. \quad \square$$

Proposition 8.2. *If $V \in C^{1,1}(\mathbb{T})$ then we have for every $\psi \in H^1(\mathbb{T})$,*

$$\lim_{\beta \rightarrow 0} \sigma_\beta^V(\psi) = \sigma_0^V(\psi).$$

Proof. Let $\psi \in H^1(\mathbb{T})$ and set $\psi_\beta := \psi - \int \psi d\mu_\beta^V$ for any $\beta \geq 0$. In particular ψ_β is continuous on \mathbb{T} and Lemma 2.3 yields

$$\lim_{\beta \rightarrow 0} \|\psi_\beta\|_{L^2(\mu_\beta^V)} = \|\psi_0\|_{L^2(\mu_0^V)} = \sigma_0^V(\psi). \tag{61}$$

We also use the integration by part formula

$$\int \phi \mathcal{A}_\beta \psi d\mu_\beta^V = \int \phi' \psi' d\mu_\beta^V \tag{62}$$

which holds for any $\phi \in H^1(\mathbb{T})$ and $\psi \in H^2(\mathbb{T})$. Note that $\psi_\beta \in H_\beta$ and, by Lemma 8.1, that $\mathcal{L}_\beta^{-1}\psi_\beta \in H^2$. Since $\mathcal{L}_\beta \geq \mathcal{A}_\beta > 0$ as operators on H_β , we obtain together with (62),

$$\begin{aligned} \sigma_\beta^V(\psi)^2 &= \left\langle \psi_\beta, \mathcal{L}_\beta^{-1}\psi_\beta \right\rangle_{H_\beta} \\ &= \left\langle \psi_\beta, \mathcal{A}_\beta \mathcal{L}_\beta^{-1}\psi_\beta \right\rangle_{L^2(\mu_\beta^V)} \\ &\leq \|\psi_\beta\|_{L^2(\mu_\beta^V)}^2. \end{aligned} \tag{63}$$

Combined with (61), this gives

$$\limsup_{\beta \rightarrow 0} \sigma_\beta^V(\psi) \leq \sigma_0^V(\psi).$$

As for the lower bound, by using (62) again, that $\mathcal{A}_\beta = \mathcal{L}_\beta - 2\pi\beta \mathcal{W}_\beta$ and $\mathcal{L}_\beta \mathcal{L}_\beta^{-1}\phi = \phi$ for every $\phi \in H^1(\mathbb{T})$, we have

$$\begin{aligned} \sigma_\beta^V(\psi)^2 &= \int \psi_\beta \mathcal{A}_\beta (\mathcal{L}_\beta^{-1}\psi_\beta) d\mu_\beta^V \\ &= \|\psi_\beta\|_{L^2(\mu_\beta^V)}^2 - 2\pi\beta \int \psi_\beta \mathcal{W}_\beta (\mathcal{L}_\beta^{-1}\psi_\beta) d\mu_\beta^V. \end{aligned} \tag{64}$$

Since $\mathcal{W}_\beta(\phi) = -H(\phi' \mu_\beta^V)$ and the Hilbert transform satisfies $H^* = -H$ on $L^2(\mathbb{T})$,

$$\int \psi_\beta \mathcal{W}_\beta (\mathcal{L}_\beta^{-1}\psi_\beta) d\mu_\beta^V = \int H(\psi_\beta \mu_\beta^V) (\mathcal{L}_\beta^{-1}\psi_\beta)' d\mu_\beta^V. \tag{65}$$

Since μ_β^V is bounded by Proposition 2.1, we have $\psi_\beta \mu_\beta^V \in L^2(\mathbb{T})$, and so does $H(\psi_\beta \mu_\beta^V)$ which moreover satisfies $\int_{\mathbb{T}} H(\psi_\beta \mu_\beta^V) dx = 0$. As a consequence, $H(\psi_\beta \mu_\beta^V)$ has a primitive $\vartheta_\beta : \mathbb{T} \rightarrow \mathbb{R}$ that we can pick so that $\int \vartheta_\beta d\mu_\beta^V = 0$. Thus, $\vartheta_\beta \in H_\beta$ and we obtain by using the Cauchy-Schwarz inequality (recalling that $\mathcal{L}_\beta^{-1} > 0$ on H_β) and (63),

$$\begin{aligned}
 \left| \int H(\psi_\beta \mu_\beta^V) (\mathcal{L}_\beta^{-1} \psi_\beta)' d\mu_\beta^V \right| &= \left| \left\langle \vartheta_\beta, \mathcal{L}_\beta^{-1} \psi_\beta \right\rangle_{\mathbb{H}_\beta} \right| \\
 &\leq \sqrt{\left\langle \vartheta_\beta, \mathcal{L}_\beta^{-1} \vartheta_\beta \right\rangle_{\mathbb{H}_\beta} \left\langle \psi_\beta, \mathcal{L}_\beta^{-1} \psi_\beta \right\rangle_{\mathbb{H}_\beta}} \\
 &\leq \|\vartheta_\beta\|_{L^2(\mu_\beta^V)} \|\psi_\beta\|_{L^2(\mu_\beta^V)}.
 \end{aligned} \tag{66}$$

To bound the term $\|\vartheta_\beta\|_{L^2(\mu_\beta^V)}$, first note that the variational constant C_β^V from (16) satisfies

$$C_\beta^V = 2F_\beta(\mu_\beta^V) - \mathcal{K}(\mu_\beta^V | \mu_0^V) \leq 2F_\beta(\mu_0^V) = 2\beta\mathcal{E}(\mu_0^V),$$

where we used that $\mathcal{K}(\mu_\beta^V | \mu_0^V) \geq 0$, that μ_β^V is the minimizer of F_β , and that $\mathcal{K}(\mu_0^V | \mu_0^V) = 0$. Thus, since $U^{\mu_\beta^V} \geq 0$, this yields together with Proposition 2.1(c) that $\mu_\beta^V(x) \leq e^{2\beta\mathcal{E}(\mu_0^V) - V(x)}$ on \mathbb{T} . In particular, there exists $C = C(V) > 0$ such that, for any $\beta \in [0, 1]$, we have $\|\mu_\beta^V\|_{L^\infty} \leq C^2/\pi$. As a consequence, using (55), we obtain for $\beta \in [0, 1]$,

$$\begin{aligned}
 \|\vartheta_\beta\|_{L^2(\mu_\beta^V)} &\leq \|\vartheta_\beta\|_{L^\infty} \\
 &\leq 2\pi \|H(\psi_\beta \mu_\beta^V)\|_{L^2} \\
 &= 2\pi \|\psi_\beta \mu_\beta^V\|_{L^2} \\
 &\leq C \|\psi_\beta\|_{L^2(\mu_\beta^V)}.
 \end{aligned}$$

Combined with (64)–(66), this finally yields

$$\liminf_{\beta \rightarrow 0} \sigma_\beta^V(\psi)^2 \geq \liminf_{\beta \rightarrow 0} \|\psi_\beta\|_{L^2(\mu_\beta^V)}^2 (1 - 2\pi C\beta) = \sigma_0^V(\psi)^2,$$

where the last identity follows from (61). The proof of the proposition is thus complete. \square

Proposition 8.3. *If $V \in C^{1,1}(\mathbb{T})$ and $\beta \min_{\mathbb{T}} \mu_\beta^V \rightarrow \infty$ as $\beta \rightarrow \infty$, then for any $\psi \in H^2(\mathbb{T})$,*

$$\lim_{\beta \rightarrow \infty} \beta \sigma_\beta^V(\psi)^2 = \|\psi\|_{H^{1/2}}^2. \tag{67}$$

If we assume instead that $\beta \min_{\mathbb{T}} \mu_\beta^{\beta V} \rightarrow \infty$ as $\beta \rightarrow \infty$, then we also have

$$\lim_{\beta \rightarrow \infty} \beta \sigma_\beta^{\beta V}(\psi)^2 = \|\psi\|_{H^{1/2}}^2. \tag{68}$$

Remark 8.1. Let us comment on the assumptions of Proposition 8.3. First, the condition that $\psi \in H^2(\mathbb{T})$ seems only technical and we expect the result still holds provided that $\psi \in H^1(\mathbb{T})$. Next, we know from Proposition 2.1(b) that $\min_{\mathbb{T}} \mu_\beta^V > 0$ and $\min_{\mathbb{T}} \mu_\beta^{\beta V} > 0$ for every fixed $\beta > 0$. However, we expect that the later quantity decays to zero as

$\beta \rightarrow \infty$. Indeed, one can verify from the Euler-Lagrange equation that if $V \in C^{1,1}(\mathbb{T})$ is not constant, then the minimizer μ_∞^V of the functional (8) does not have full support on \mathbb{T} . On the other-hand, if the potential V is fixed, then we already know from Lemma 2.3 that $\mu_\beta^V \rightarrow \frac{dx}{2\pi}$ weakly. In Lemma 8.4 below, we establish that, if this convergence holds in L^p for $p > 1$ with a rate of at most $c \log \beta / \beta$ for $c > 0$ small enough, then the hypothesis that $\beta \min_{\mathbb{T}} \mu_\beta^V \rightarrow +\infty$ as $\beta \rightarrow +\infty$ is satisfied.

We are now ready to prove of Proposition 8.3.

Proof. We start by proving (67). Recall that by definition, we have for every $\phi \in H^1(\mathbb{T})$,

$$\|\phi\|_{H^{1/2}}^2 = \langle \phi', H(\phi) \rangle_{L^2} = -\langle H(\phi'), \phi \rangle_{L^2}.$$

By Lemma 6.1, the operator \mathcal{W}_β^{-1} is well-defined on H_β . Moreover, by (49) and since $H^{-1} = -H$ on $L_0^2(\mathbb{T})$, we have for every $\phi \in H_\beta$,

$$(\mathcal{W}_\beta^{-1}\phi)' \mu_\beta^V = -H\phi + \int_{\mathbb{T}} (\mathcal{W}_\beta^{-1}\phi)' d\mu_\beta^V.$$

Recall that $\psi_\beta = \psi - \int \psi d\mu_\beta^V$ for any $\beta \geq 0$ and that $\psi_\beta \in H_\beta$. Using further that $\mathcal{L}_\beta \geq 2\pi\beta\mathcal{W}_\beta$ as operators on H_β , we obtain for every $\beta > 0$ the upper bound,

$$\begin{aligned} \beta\sigma_\beta^V(\psi)^2 &= \beta \left\langle \psi_\beta, \mathcal{L}_\beta^{-1}\psi_\beta \right\rangle_{H_\beta} \\ &\leq \frac{1}{2\pi} \left\langle \psi_\beta, \mathcal{W}_\beta^{-1}\psi_\beta \right\rangle_{H_\beta} \\ &= -\left\langle \psi'_\beta, H(\psi_\beta) \right\rangle_{L^2} \\ &= \|\psi_\beta\|_{H^{1/2}}^2 \\ &= \|\psi\|_{H^{1/2}}^2. \end{aligned} \tag{69}$$

As for the lower bound, recalling that $\mathcal{W}_\beta(\phi) = -H(\phi' \mu_\beta^V)$ and writing $2\pi\beta\mathcal{W}_\beta = \mathcal{L}_\beta - \mathcal{A}_\beta$, since H is an isometry of $L_0^2(\mathbb{T})$ and $\widehat{\psi}'_0 = 0$, we obtain

$$\begin{aligned} \beta\sigma_\beta^V(\psi)^2 &= 2\pi\beta \left\langle \psi', (\mathcal{L}_\beta^{-1}\psi_\beta)' \mu_\beta^V \right\rangle_{L^2} \\ &= 2\pi\beta \left\langle H(\psi'), H((\mathcal{L}_\beta^{-1}\psi_\beta)' \mu_\beta^V) \right\rangle_{L^2} \\ &= -2\pi\beta \left\langle H(\psi'), \mathcal{W}_\beta(\mathcal{L}_\beta^{-1}\psi_\beta) \right\rangle_{L^2} \\ &= -\langle H(\psi'), \psi_\beta \rangle_{L^2} + \left\langle H(\psi'), \mathcal{A}_\beta(\mathcal{L}_\beta^{-1}\psi_\beta) \right\rangle_{L^2} \\ &= \|\psi\|_{H^{1/2}}^2 + \left\langle H(\psi'), \mathcal{A}_\beta(\mathcal{L}_\beta^{-1}\psi_\beta) \right\rangle_{L^2}. \end{aligned} \tag{70}$$

Now, we set

$$\vartheta_\beta := \frac{H(\psi')}{\mu_\beta^V}.$$

Since by assumption $\psi'' \in L^2(\mathbb{T})$, $\|(\mu_\beta^V)'\|_{L^\infty} < \infty$ by Proposition 2.5 and H maps $L^2(\mathbb{T})$ into $L^2_0(\mathbb{T})$, it easily follows from Proposition 2.1(b) that $\vartheta_\beta \in \mathbf{H}_\beta$. Moreover, since $\mathcal{L}_\beta^{-1}(\psi_\beta) \in \mathbf{H}^2(\mathbb{T})$ according to Lemma 8.1, we can use formula (62) and the Cauchy-Schwarz inequality (recalling that $\mathcal{L}_\beta^{-1} > 0$ on \mathbf{H}_β) to obtain

$$\begin{aligned} \left\langle H(\psi'), \mathcal{A}_\beta(\mathcal{L}_\beta^{-1}\psi_\beta) \right\rangle_{L^2} &= \frac{1}{2\pi} \left\langle \vartheta_\beta, \mathcal{A}_\beta(\mathcal{L}_\beta^{-1}\psi_\beta) \right\rangle_{L^2(\mu_\beta^V)} \\ &= -\frac{1}{2\pi} \left\langle \vartheta_\beta, \mathcal{L}_\beta^{-1}\psi_\beta \right\rangle_{\mathbf{H}_\beta} \\ &\geq -\frac{1}{2\pi} \sigma_\beta^V(\psi) \sigma_\beta^V(\vartheta_\beta). \end{aligned} \tag{71}$$

Next, using (63) and Proposition 2.1(b), we then have

$$\frac{1}{2\pi} \sigma_\beta^V(\vartheta_\beta) \leq \|\vartheta_\beta\|_{L^2(\mu_\beta^V)} \leq \delta^{-1/2} \|H(\psi')\|_{L^2} = \delta^{-1/2} \|\psi\|_{\mathbf{H}^1}$$

for some $\delta = \delta(\beta) > 0$ that satisfies, by assumption, $\beta\delta \rightarrow \infty$ as $\beta \rightarrow \infty$. Combined with (70)–(71) this yields

$$\beta \sigma_\beta^V(\psi)^2 \geq \|\psi\|_{\mathbf{H}^{1/2}}^2 - \delta^{-1/2} \|\psi\|_{\mathbf{H}^1} \sigma_\beta^V(\psi).$$

By computing the roots of the polynomial function $x \mapsto x^2 - (\beta\delta)^{-1/2} \|\psi\|_{\mathbf{H}^1} x + \|\psi\|_{\mathbf{H}^{1/2}}^2$ this provides in turn,

$$\liminf_{\beta \rightarrow +\infty} \sqrt{\beta} \sigma_\beta^V(\psi) \geq \liminf_{\beta \rightarrow +\infty} \sqrt{\|\psi\|_{\mathbf{H}^{1/2}}^2 + \frac{\|\psi\|_{\mathbf{H}^1}^2}{4\beta\delta} - \frac{\|\psi\|_{\mathbf{H}^1}}{2\sqrt{\beta\delta}}} = \|\psi\|_{\mathbf{H}^{1/2}}$$

and, together with the upper bound (69), the claim (67) is proven.

Since the proof of (68) is identical to the one of (67) after replacing μ_β^V by $\mu_\beta^{\beta V}$ everywhere in the above arguments, the proposition is obtained. \square

Lemma 8.4. *Let $V \in \mathbf{H}^1(\mathbb{T})$ and $p > 1$. Suppose that for all β sufficiently large,*

$$\|\mu_\beta^V - \frac{dx}{2\pi}\|_{L^p} \leq \kappa_p \frac{\log \beta}{\beta}$$

for a constant $\kappa_p > 0$ which is sufficiently small, then $\beta \inf_{\mathbb{T}} \mu_\beta^V \rightarrow +\infty$ as $\beta \rightarrow +\infty$.

Proof. Recall that $U^{\frac{dx}{2\pi}} = \log 2$. Since $\mathcal{K}(\mu_\beta^V | \mu_0^V) \geq 0$ and $\frac{dx}{2\pi}$ minimizes \mathcal{E} , we have for $\beta > 0$,

$$C_\beta^V \geq 2\beta\mathcal{E}(\mu_\beta^V) \geq 2\beta \log 2.$$

From (16), which holds for all $x \in \mathbb{T}$ if $V \in H^1(\mathbb{T})$ according to Proposition 2.5, the density μ_β^V satisfies for all $x \in \mathbb{T}$,

$$\mu_\beta^V(x) = e^{C_\beta^V - 2\beta U^{\mu_\beta^V}(x) - V(x)} \geq e^{2\beta(U^{\frac{dx}{2\pi}} - U^{\mu_\beta^V}(x))} \inf_{\mathbb{T}} e^{-V}.$$

Next, notice that the function $g = \log \left| \sin \left(\frac{\cdot}{2} \right) \right|^{-1}$ is L^p for any $p > 0$. Using Young's convolution inequality we obtain, for every $p > 1$,

$$\|U^{\mu_\beta^V} - U^{\frac{dx}{2\pi}}\|_{L^\infty} = \|g * (\mu_\beta^V - \frac{1}{2\pi})\|_{L^\infty} \leq \|g\|_{L^{\frac{p}{p-1}}} \|\mu_\beta^V - \frac{1}{2\pi}\|_{L^p}$$

and the lemma follows. \square

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