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## **Three-Dimensional Spin Rotations at the Fermi Surface of a Strongly Spin-Orbit Coupled Surface System**

Höpfner, P ; Schäfer, J ; Fleszar, A ; Dil, J H ; Slomski, B ; Meier, F ; Loho, C ; Blumenstein, C ; Patthey, L ; Hanke, W ; Claessen, R

**Abstract:** The spin texture of the metallic two-dimensional electron system ( $\sqrt{3}\times\sqrt{3}$ )-Au/Ge(111) is revealed by fully three-dimensional spin-resolved photoemission, as well as by density functional calculations. The large hexagonal Fermi surface, generated by the Au atoms, shows a significant splitting due to spin-orbit interactions. The planar components of the spin exhibit a helical character, accompanied by a strong out-of-plane spin component with alternating signs along the six Fermi surface sections. Moreover, in-plane spin rotations toward a radial direction are observed close to the hexagon corners. Such a threefold-symmetric spin pattern is not described by the conventional Rashba model. Instead, it reveals an interplay with Dresselhaus-like spin-orbit effects as a result of the crystalline anisotropies.

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Supplemental Material

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# Supplemental Material

## Three-Dimensional Spin Rotations at the Fermi Surface of a Strongly Spin-Orbit Coupled Surface System

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L. Patthey, W. Hanke, and R. Claessen

This supplemental material provides details of the band-structure modeling according to Eq. (3):

$$H(k) = \left( \frac{\hbar^2 k^2}{2m^*} - C + c_h(k_+^6 + k_-^6) \right) \sigma_0 + v(k_x \sigma_y - k_y \sigma_x) + \lambda(k_+^3 + k_-^3) \sigma_z + i\zeta(k_+^5 \sigma_+ - k_-^5 \sigma_-),$$

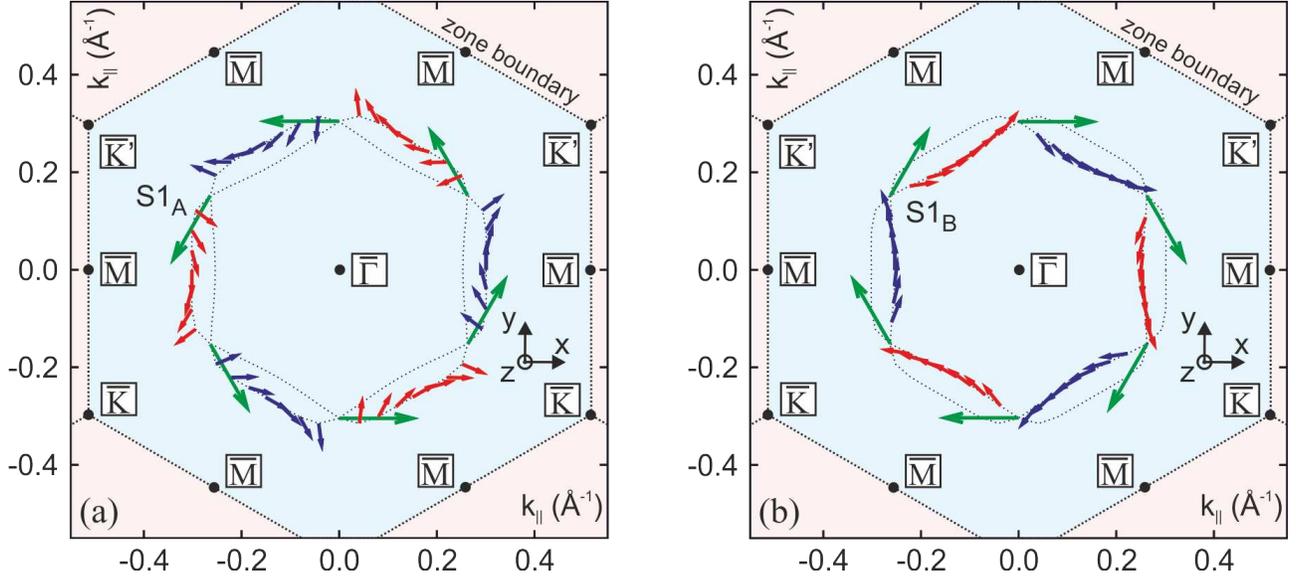
where  $\sigma_0$  is the  $2 \times 2$  identity matrix,  $\sigma_x$ ,  $\sigma_y$  and  $\sigma_z$  are Pauli matrices,  $\sigma_{\pm} = \sigma_x \pm i\sigma_y$ , and  $k_{\pm} = k_x \pm ik_y$ . This Hamiltonian is invariant with respect to symmetry operations of the  $C_{3v}$  group and time-reversal symmetry. For the following choice of parameters, the  $E = 0$  energy surface of this model Hamiltonian reproduces well both, the Fermi contour and the spin structure of the metallic surface band in  $(\sqrt{3} \times \sqrt{3})$ -Au/Ge(111):  $m^* = 0.4 m_e$ ,  $C = 0.77 \text{ eV}$ ,  $c_h = 70 \text{ eV} \text{ \AA}^{-6}$ ,  $v = -0.1 \text{ eV} \text{ \AA}^{-1}$ ,  $\lambda = -7 \text{ eV} \text{ \AA}^{-3}$ ,  $\zeta = 5.5 \text{ eV} \text{ \AA}^{-5}$ .

The first part of the Hamiltonian, proportional to  $\sigma_0$ , is the spinless part, giving rise to spin doubly degenerate states. The sixth-order in  $k$  term in the parentheses is responsible for the hexagonal warping of the constant energy contours. This term reproduces well the Fermi contour of our *ab-initio* calculations, before including the spin-orbit interaction.

The second part of the Hamiltonian,  $v(k_x \sigma_y - k_y \sigma_x)$ , is the well-known Rashba term. It has full  $O(2)$  symmetry and produces an isotropic spin splitting. The two resulting constant-energy contours (corresponding to  $S1_A$  and  $S1_B$ ) have a spin texture being fully planar, vortical (clockwise for  $S1_B$  and counter-clockwise for  $S1_A$ ) and perpendicular to the surface momentum vector  $(k_x, k_y)$ .

The third part of the Hamiltonian, of third-order in  $k$ , was proposed by L. Fu in Ref. 24 in order to explain the hexagonal warping of the Dirac-fermion band at the surface of the topological insulator  $\text{Bi}_2\text{Te}_3$ . Apart from affecting the shape of energy contours, it generates a  $z$ -component in the spin vector. This vertical part of the spin exhibits an alternating sign (*up* or *down*) between neighboring hexagonal sheets of the energy contours. The in-plane part of the spin vector remains, however, perpendicular to the momentum vector  $(k_x, k_y)$ . Along the  $\bar{\Gamma} - \bar{K}$  azimuth, there is no  $z$ -component of spin due to symmetry reasons.

The last term, of fifth-order in  $k$ ,  $i\zeta(k_+^5 \sigma_+ - k_-^5 \sigma_-)$ , was recently proposed by S. Basak *et al.* [25]. This term is needed to explain the *ab-initio* results showing radial rotations of the in-plane spin components for the Dirac-fermion bands at the surface of the topological insulator  $\text{Bi}_2\text{Te}_3$ . Similarly as the third-order term, it contributes to the hexagonal warping of the energy contours, but its main function is the rotation of in-plane spin components towards a radial direction. This happens in close vicinity to the  $\bar{\Gamma} - \bar{K}$  azimuth.



**Figure S1:**  $E = 0$  energy contours derived from the model Hamiltonian with the in-plane spin texture of the spin-split states (a)  $S1_A$  and (b)  $S1_B$ . The spin arrows are red for spins directed out-of-plane and blue for an orientation into the plane. Green color signifies fully in-plane spin alignment.

Figure S1 displays the energy contours and spin texture of the model Hamiltonian (Eq. (3)) for the choice of parameters listed above. For better clarity each spin-split state ( $S1_A$  and  $S1_B$ ) is shown in a separate plot.